

# Binary Black Holes mergers from Population III stars

Filippo Santoliquido

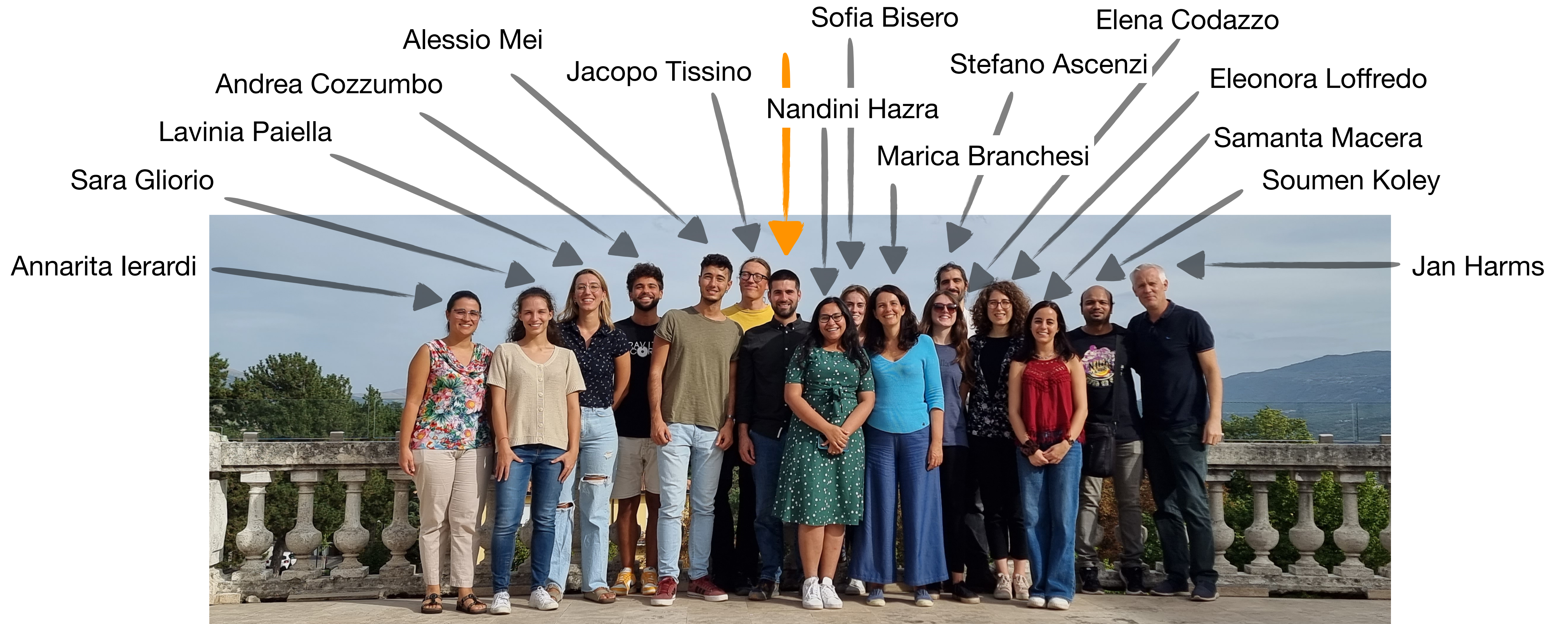
**Collaborators:** Marica Branchesi, Jan Harms, Ulyana Dupletsa, Jacopo Tissino, Giuliano Iorio, Michela Mapelli, M. Celeste Artale *et al.*

*Universidad Andrés Bello – August 6, 2024*





# GSSI team

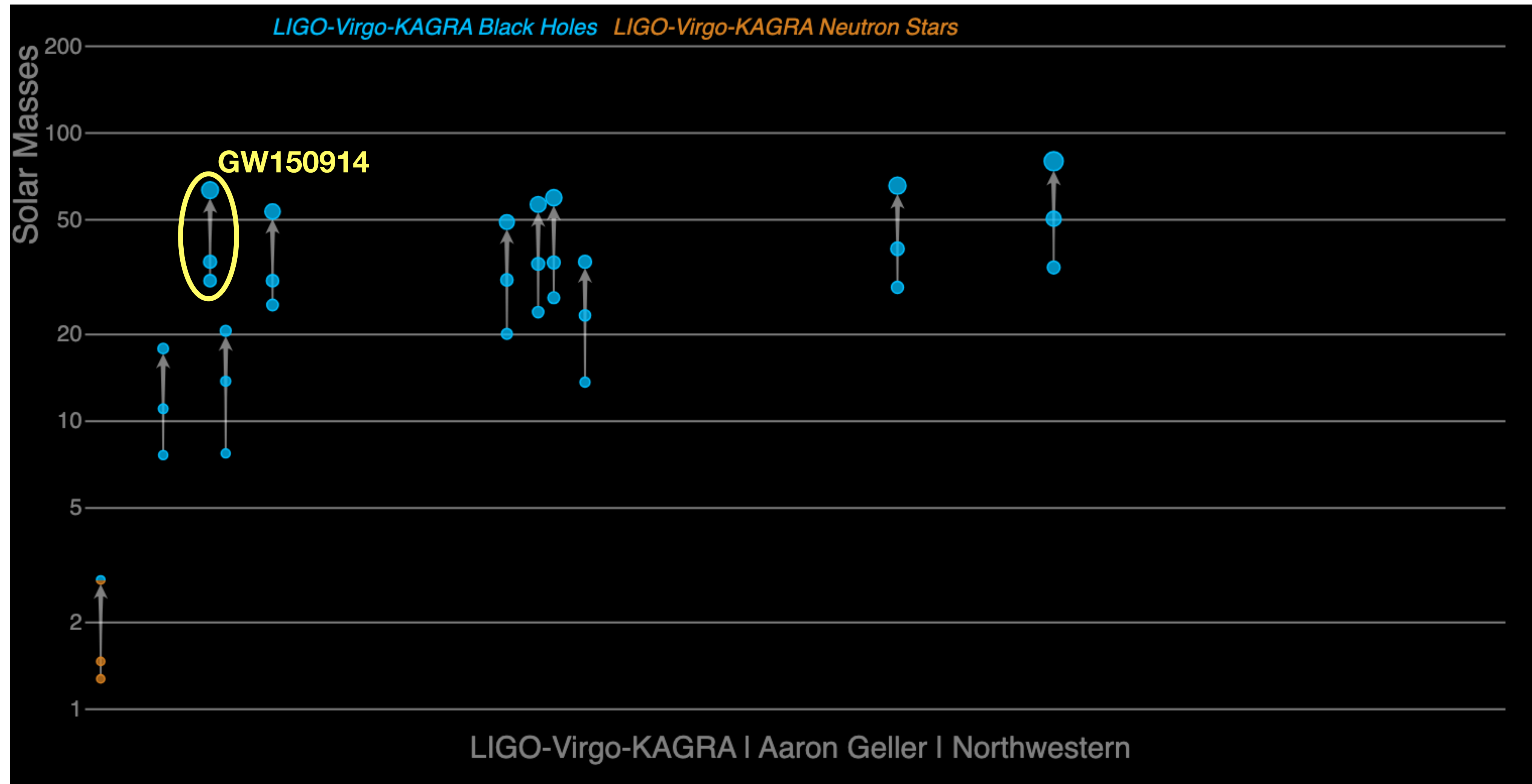


And many others...

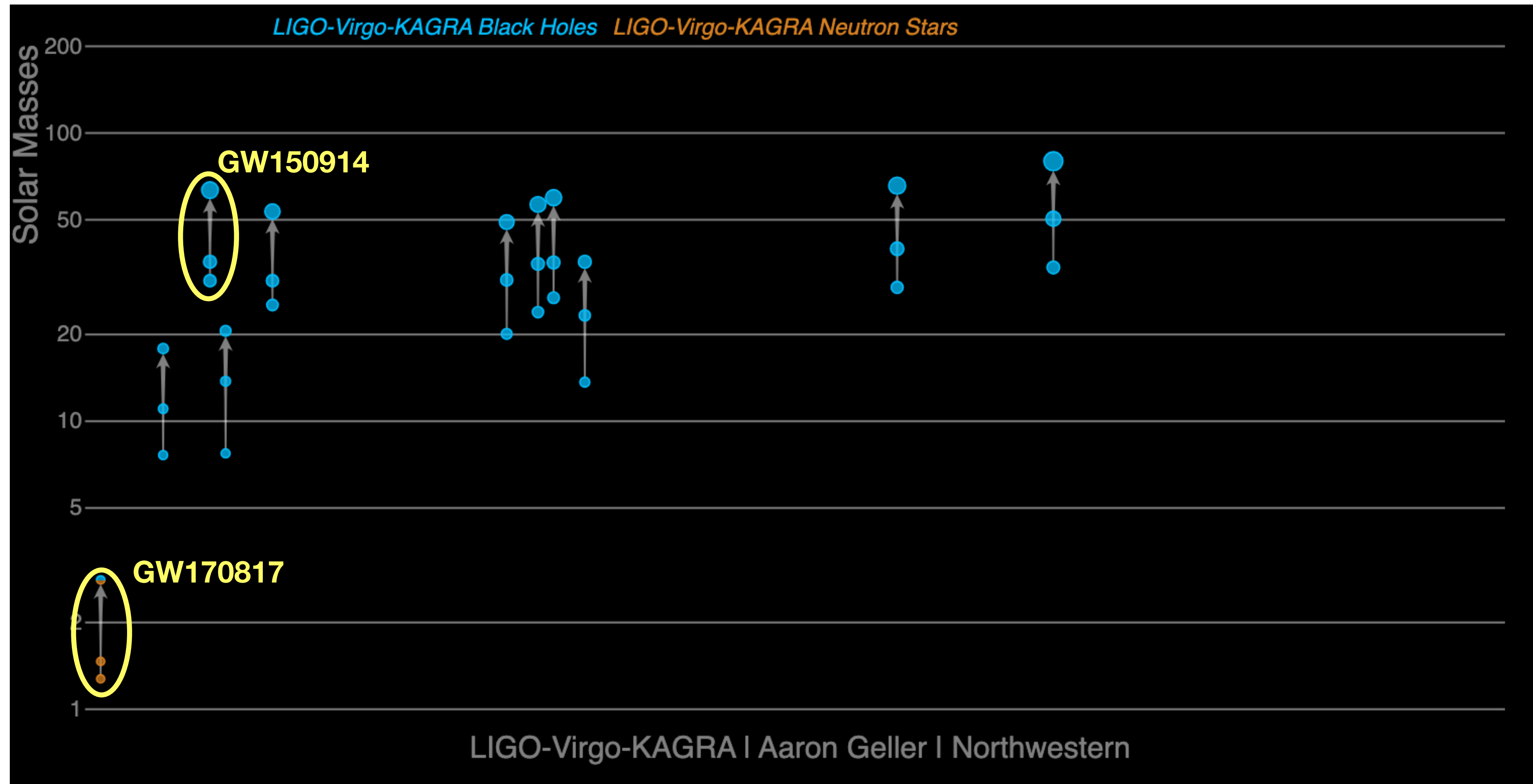
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# Gravitational Wave Astrophysics

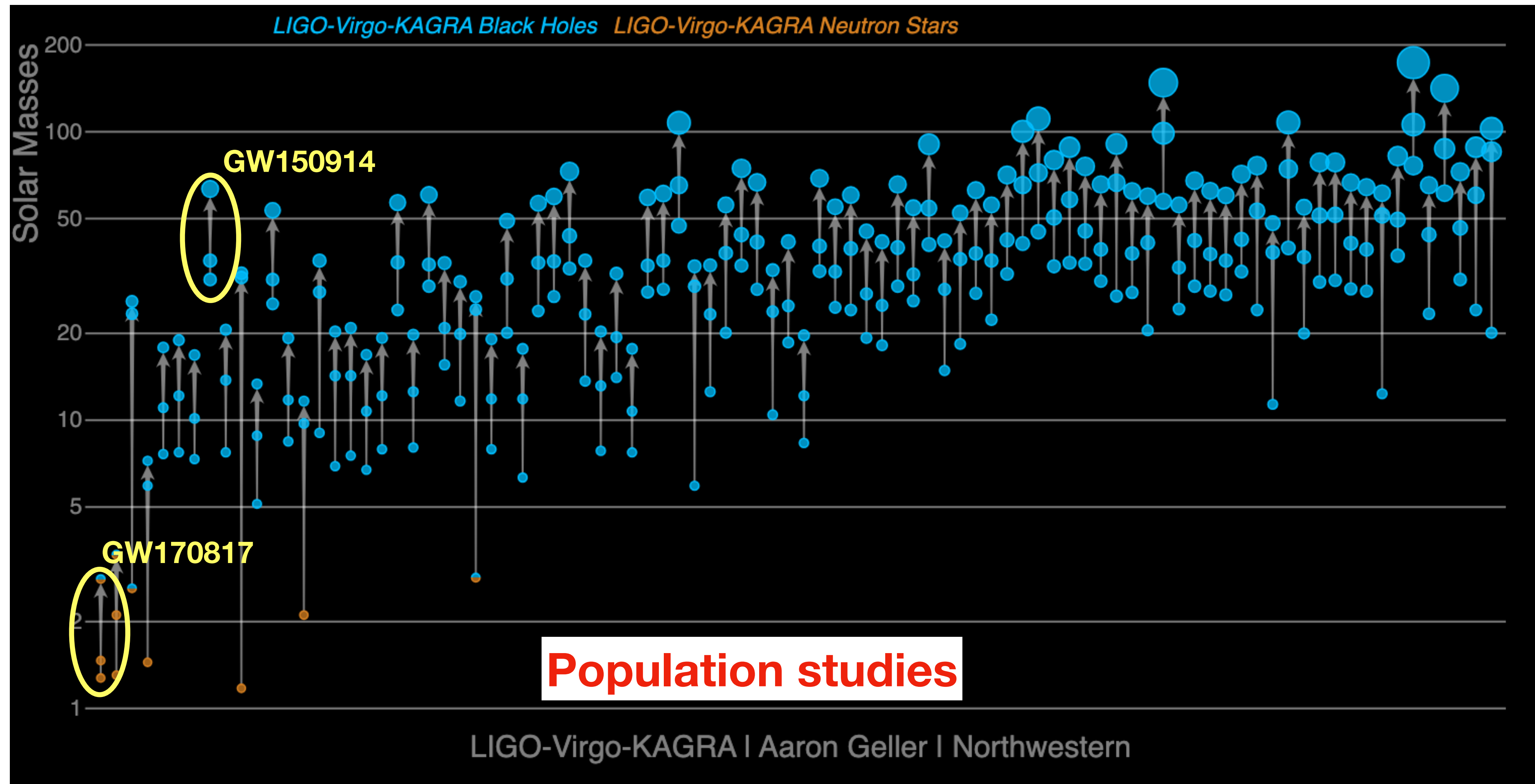


# Gravitational Wave Astrophysics

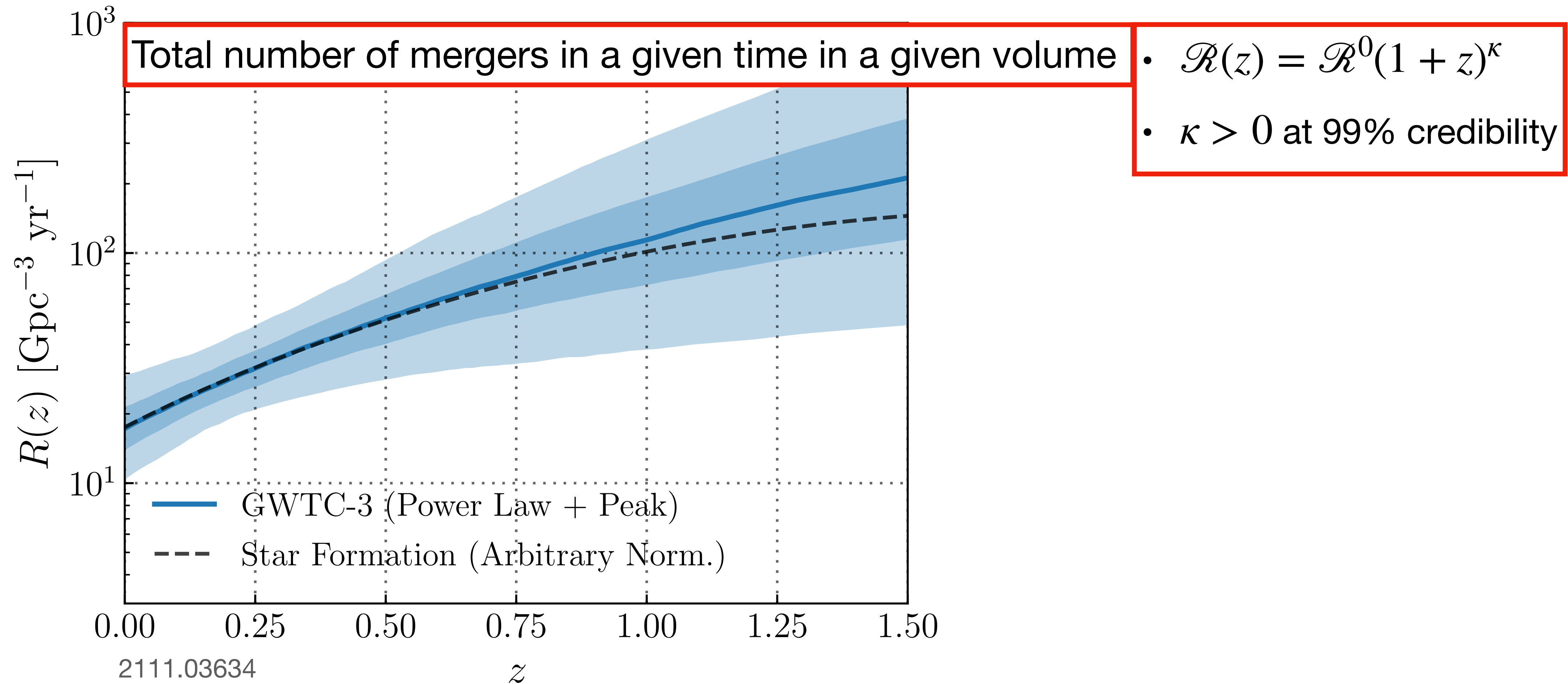




# Gravitational Wave Astrophysics

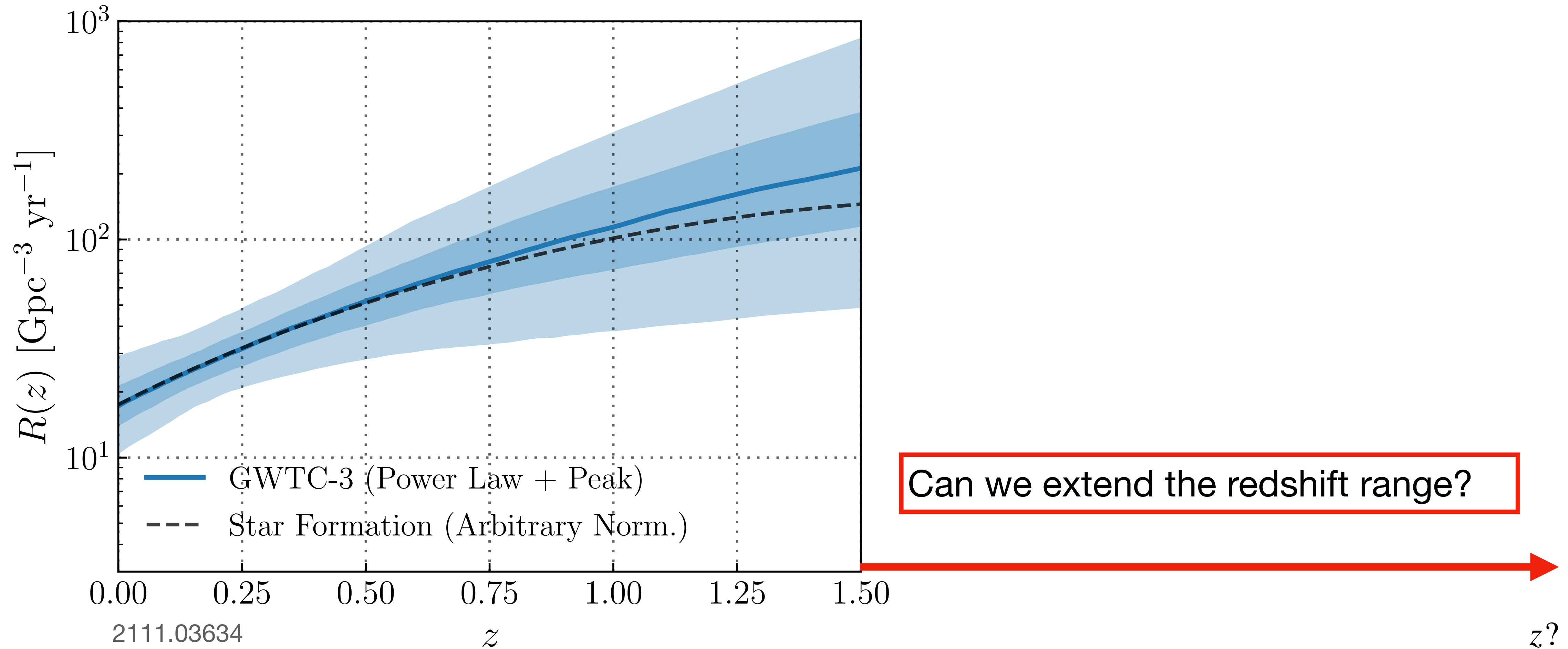


# Merger rate density



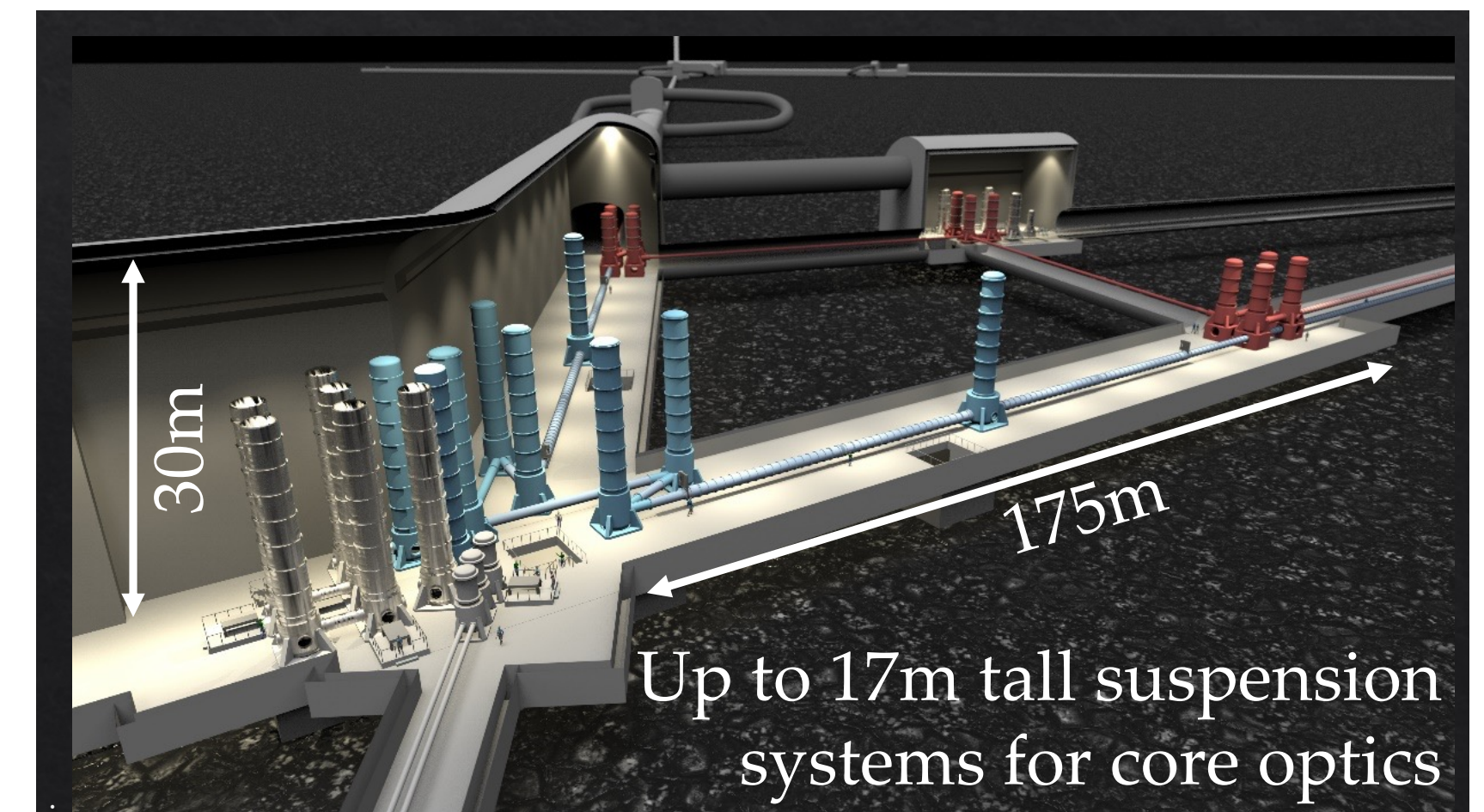
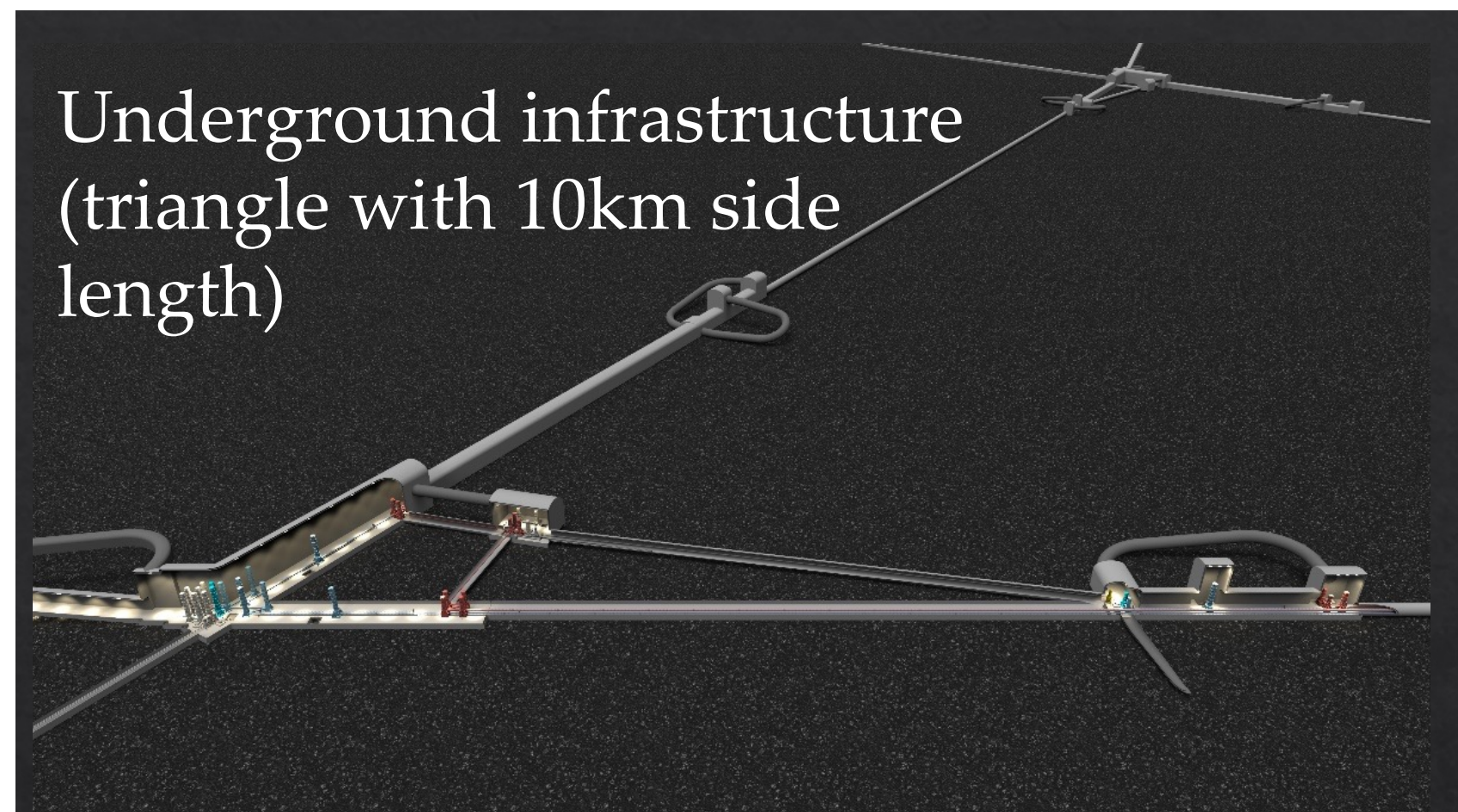


# Merger rate density





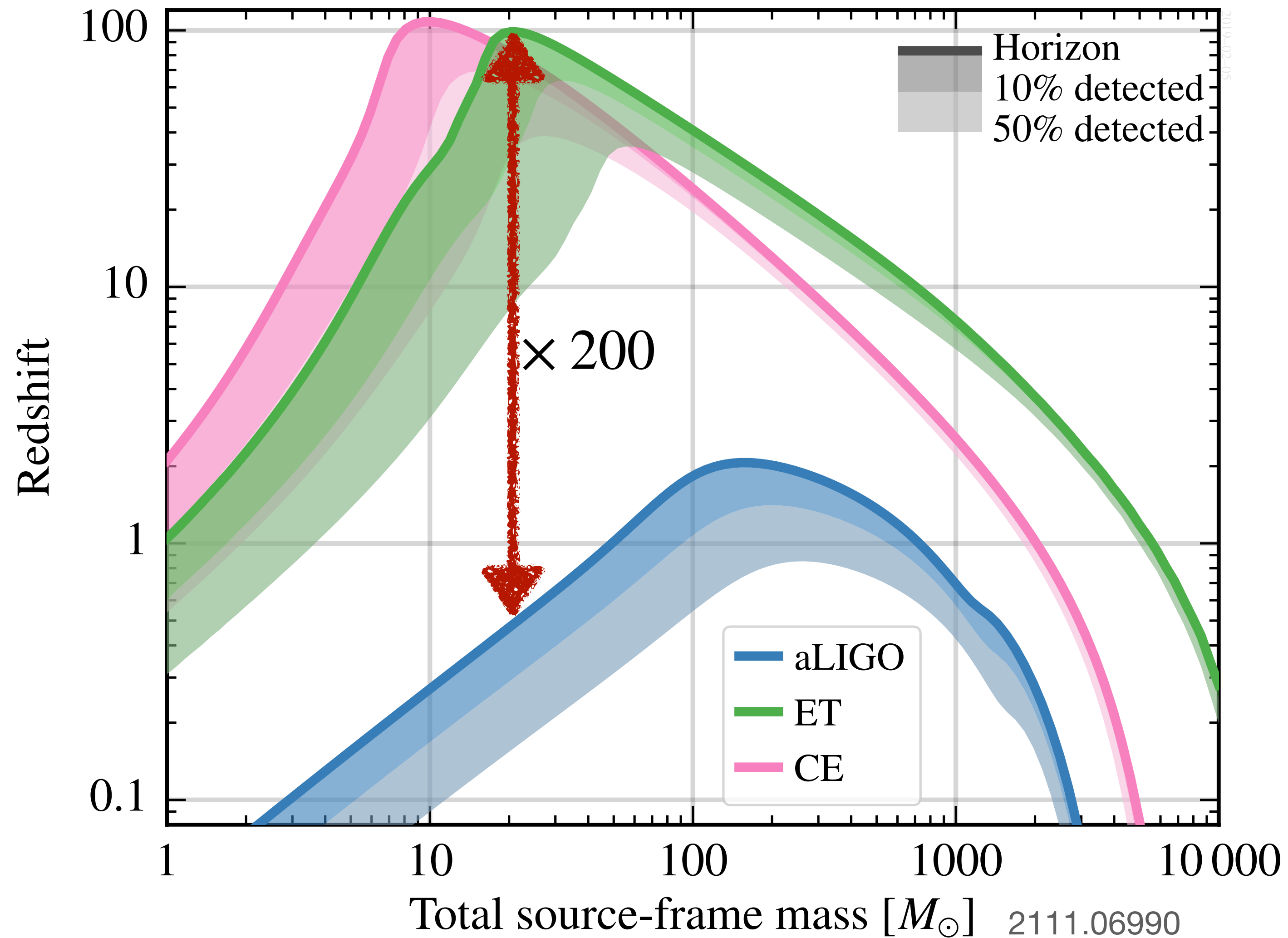
# The Einstein Telescope



Credits: Jan Harms



# Third-generation detectors



# Population III stars

- Population III stars are believed to be the first generation of stars formed at **high redshift** ( $z > 20$ )
- They are massive and formed from pristine gas (i.e. **zero metallicity**)
- They are still undetected (*traces of their existence with JWST*)
- Modelled through cosmological simulations

Ref. [Klessen & Glover 2023](#), [Bromm & Larson 2004](#), [Zackrisson et al. 2023](#), [Maiolino et al. 2024](#)

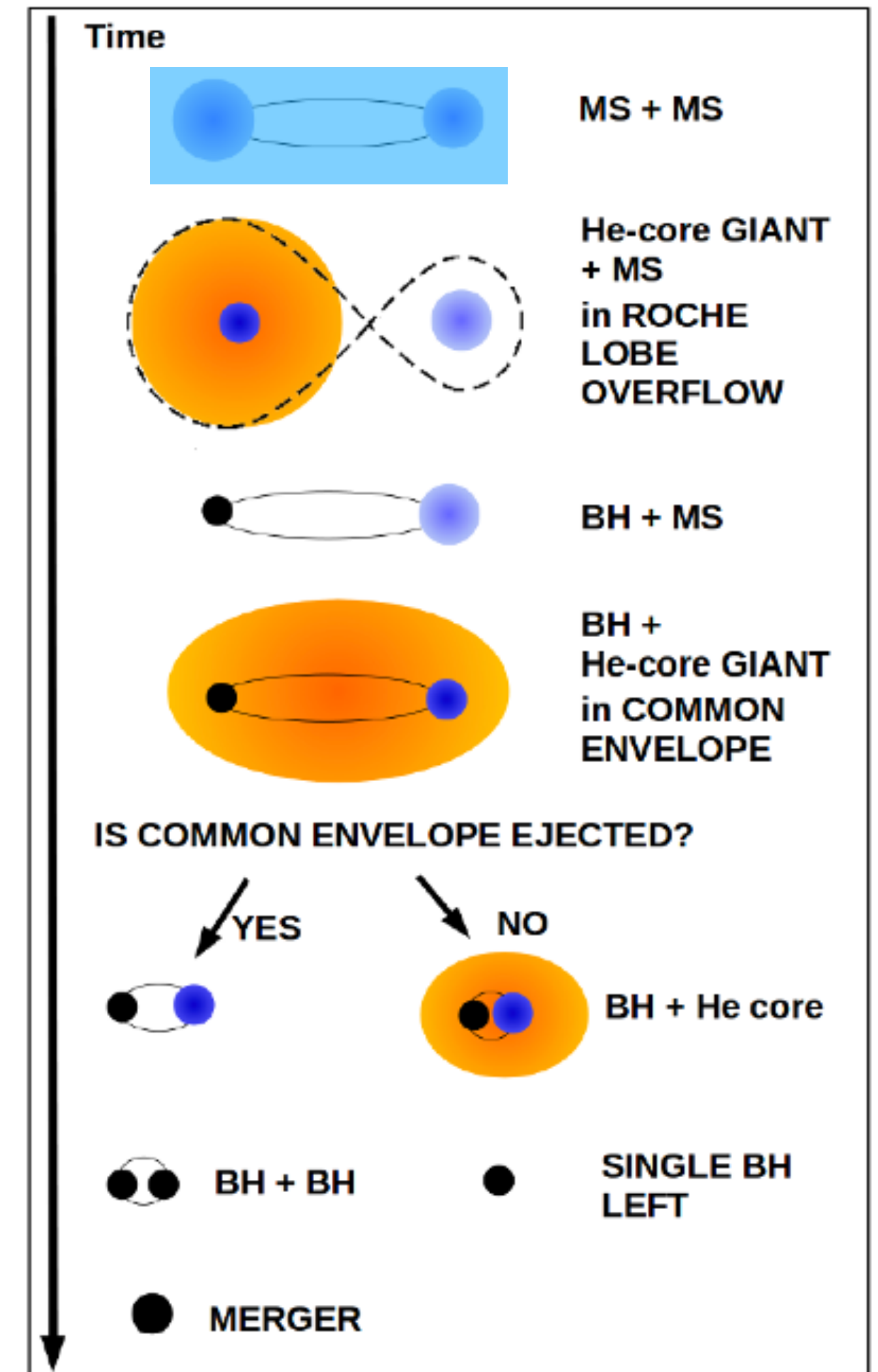


# Our goals

- Large parameter space exploration of Pop. III BBHs
- Merger rate density
- Evolution of mass spectrum with redshift
- Expected detection rates with the Einstein Telescope

# Isolated formation channel

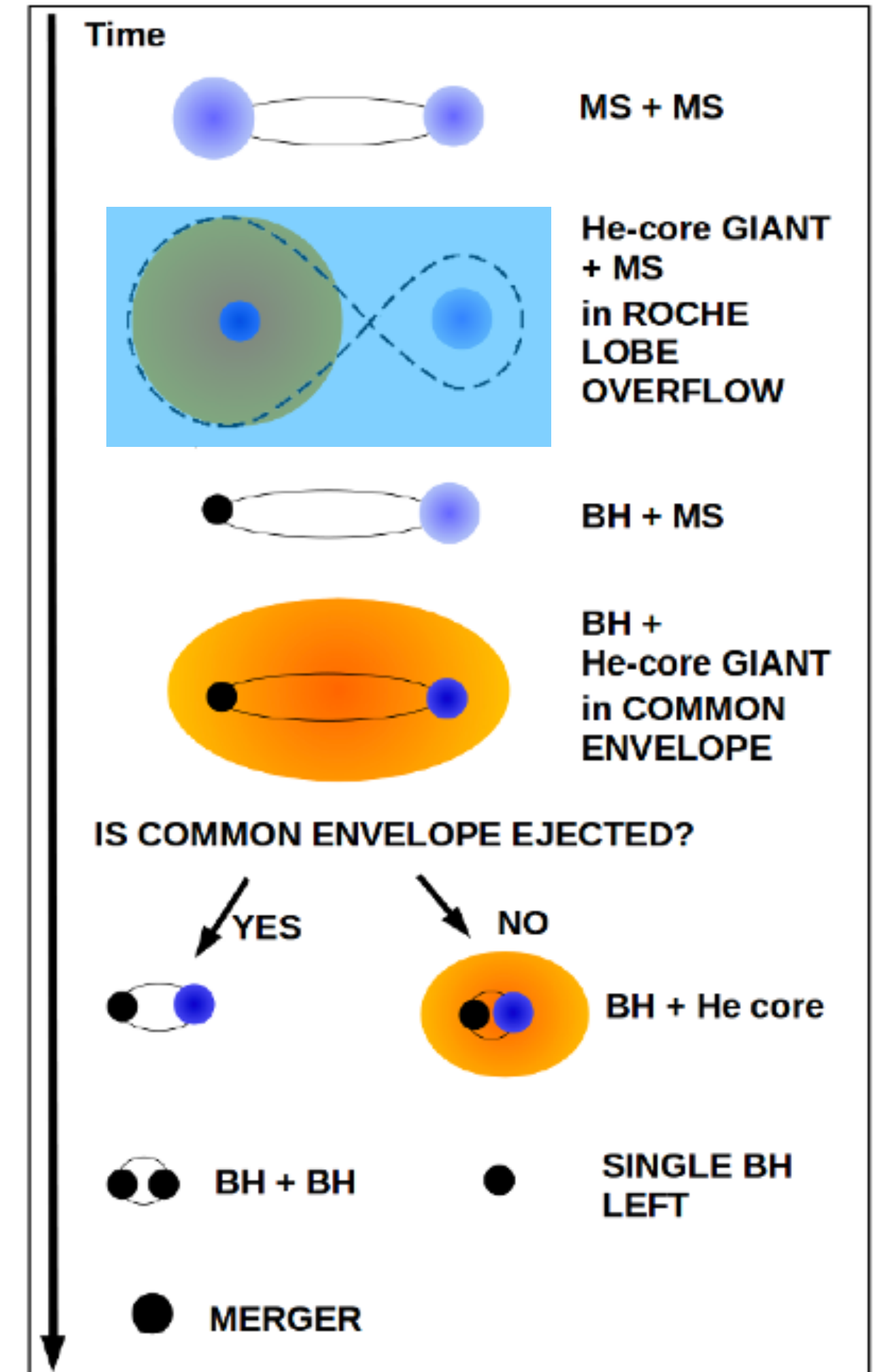
- Two massive stars form in the **same binary system** and evolve together
- *Let's see an example*





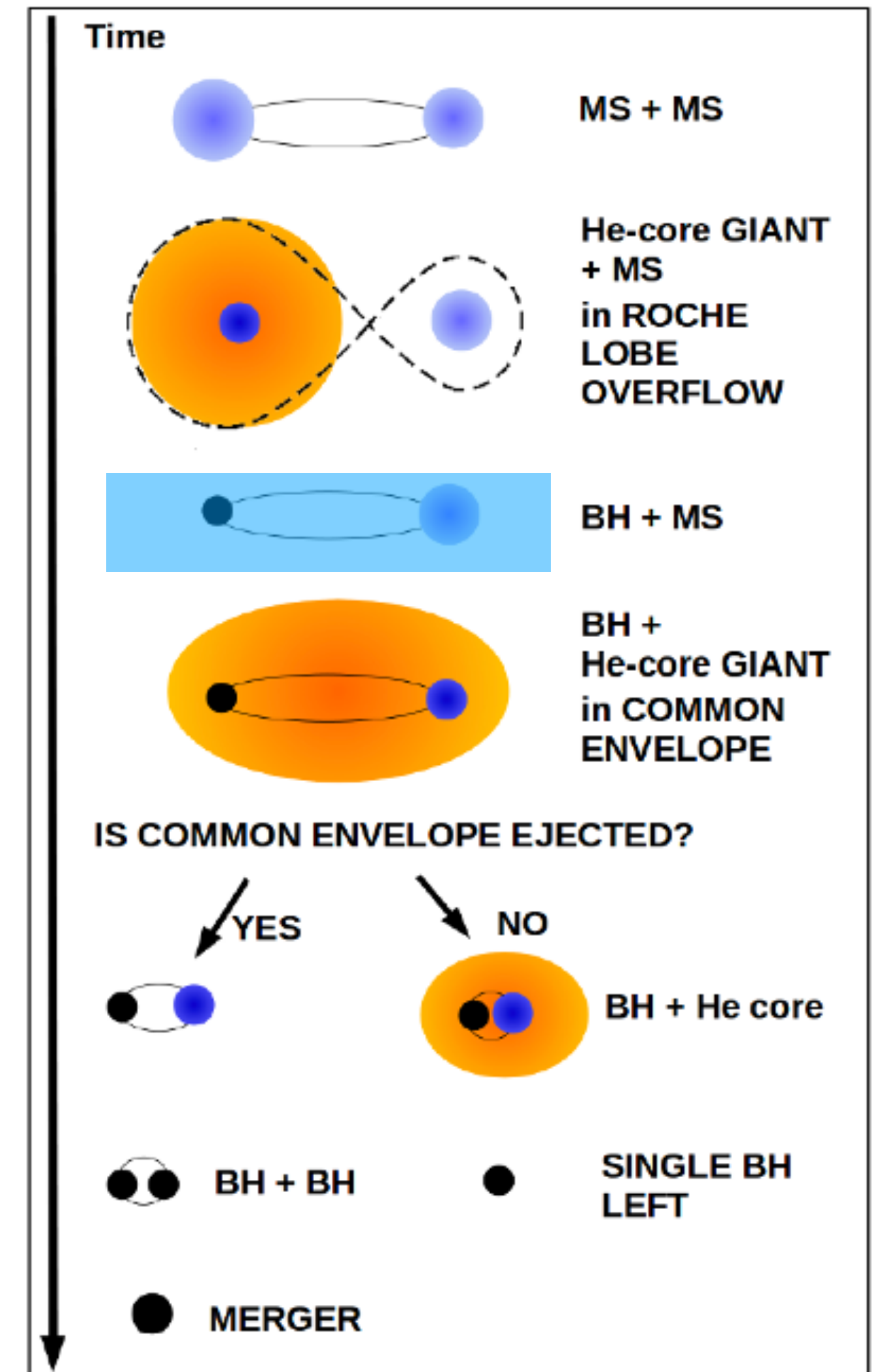
# Isolated formation channel

- Following radius expansion, one star can fill its Roche lobe
- Mass can be transferred to the other in a **stable mass transfer** (Roche lobe overflow)
- The orbit of the binary system shrinks due to angular momentum loss



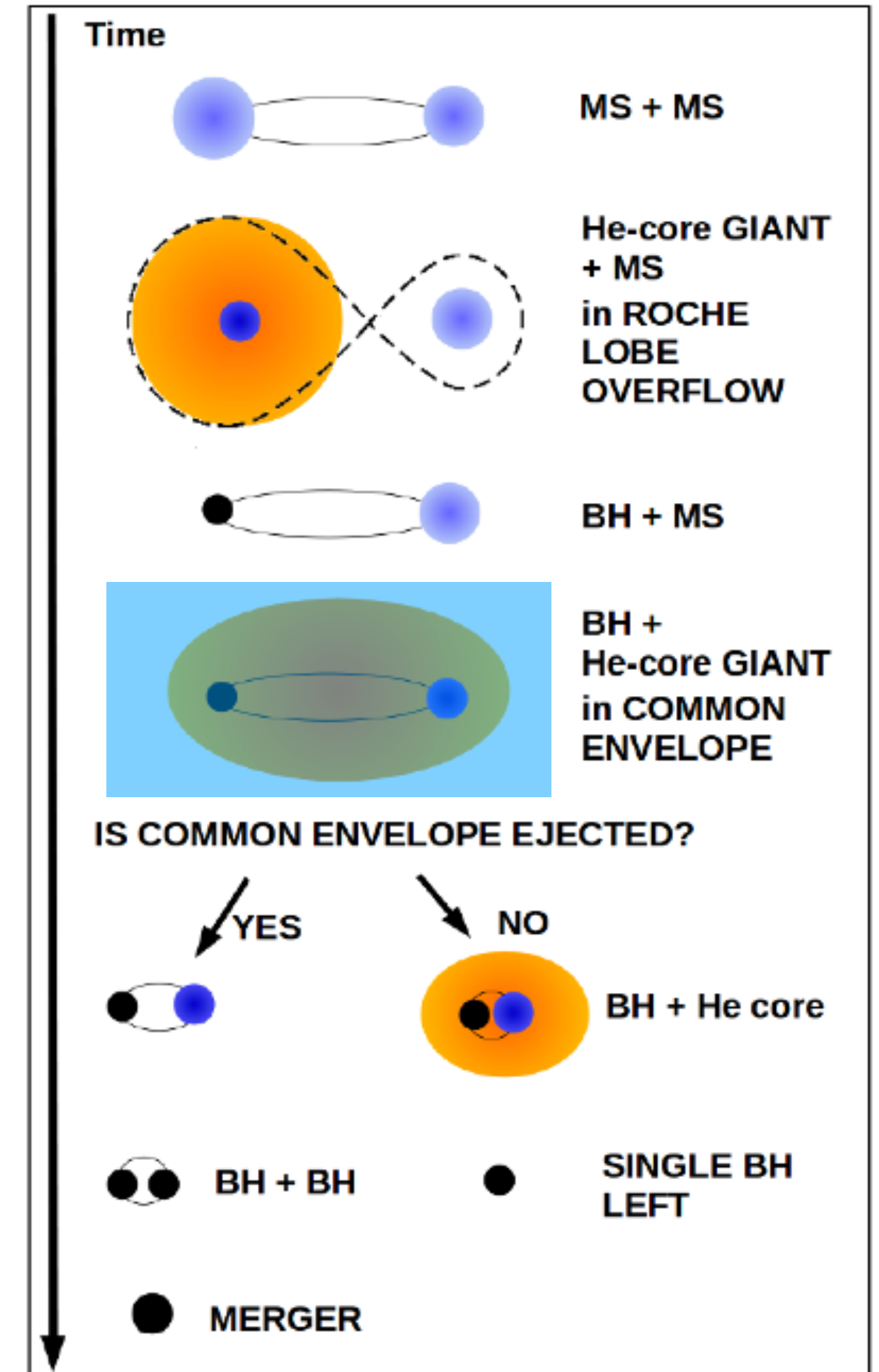
# Isolated formation channel

- If  $M_{\text{ZAMS}} \gtrsim 20 M_{\odot}$ , the first black is formed after core collapse supernova



# Isolated formation channel

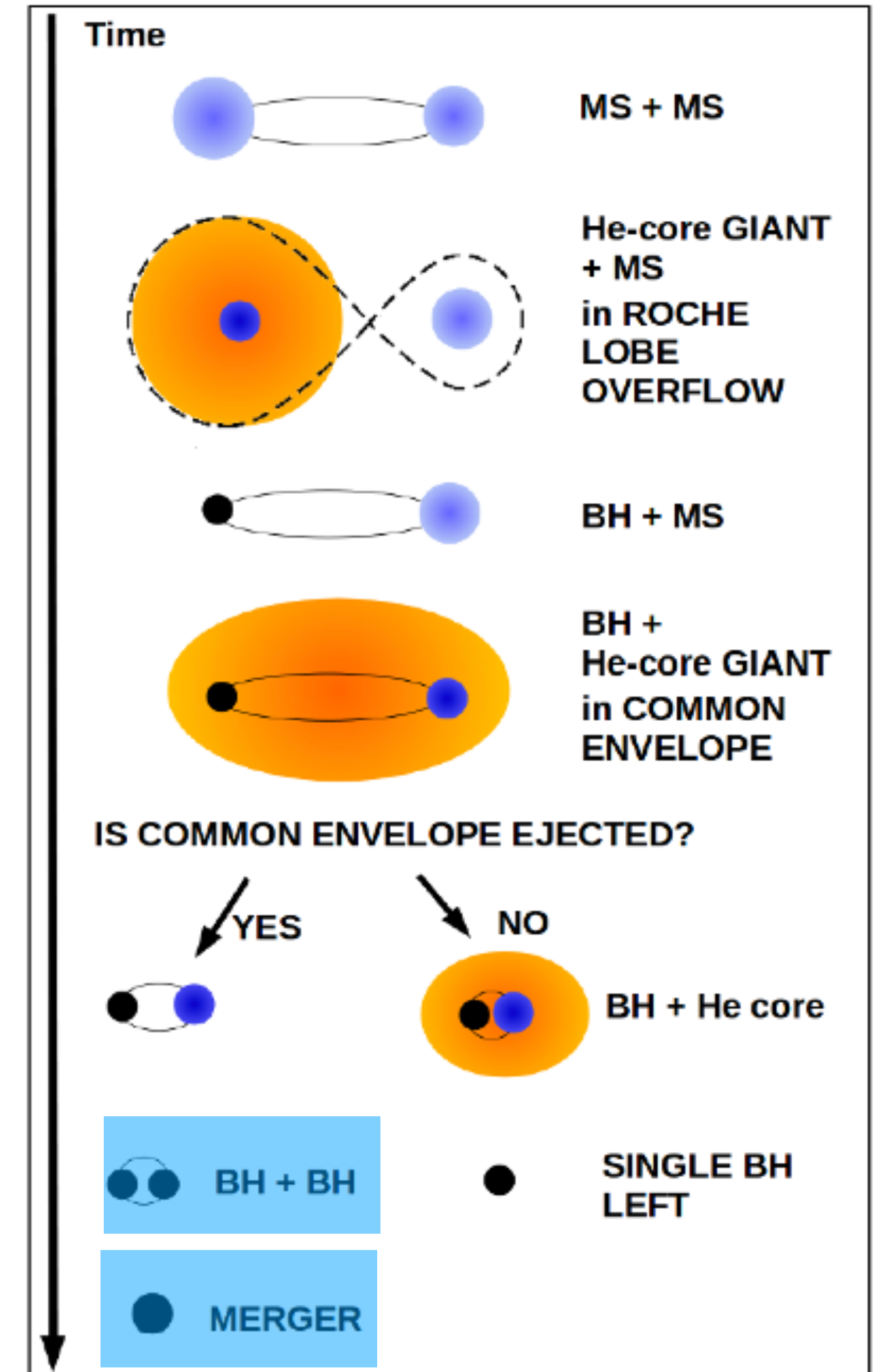
- The other star starts transferring mass through unstable mass transfer
- Beginning of the **common envelope** phase
- A drag force is exerted between the black hole and the stellar core.
- Orbital energy is transferred to the common envelope that is eventually ejected, shrinking the binary system





# Isolated formation channel

- If  $M_{\text{ZAMS}} \gtrsim 20 M_{\odot}$ , the second black hole forms
- The binary system starts to lose energy through **gravitational wave radiation emission**
- The **two black holes** might eventually **merge** within an Hubble time

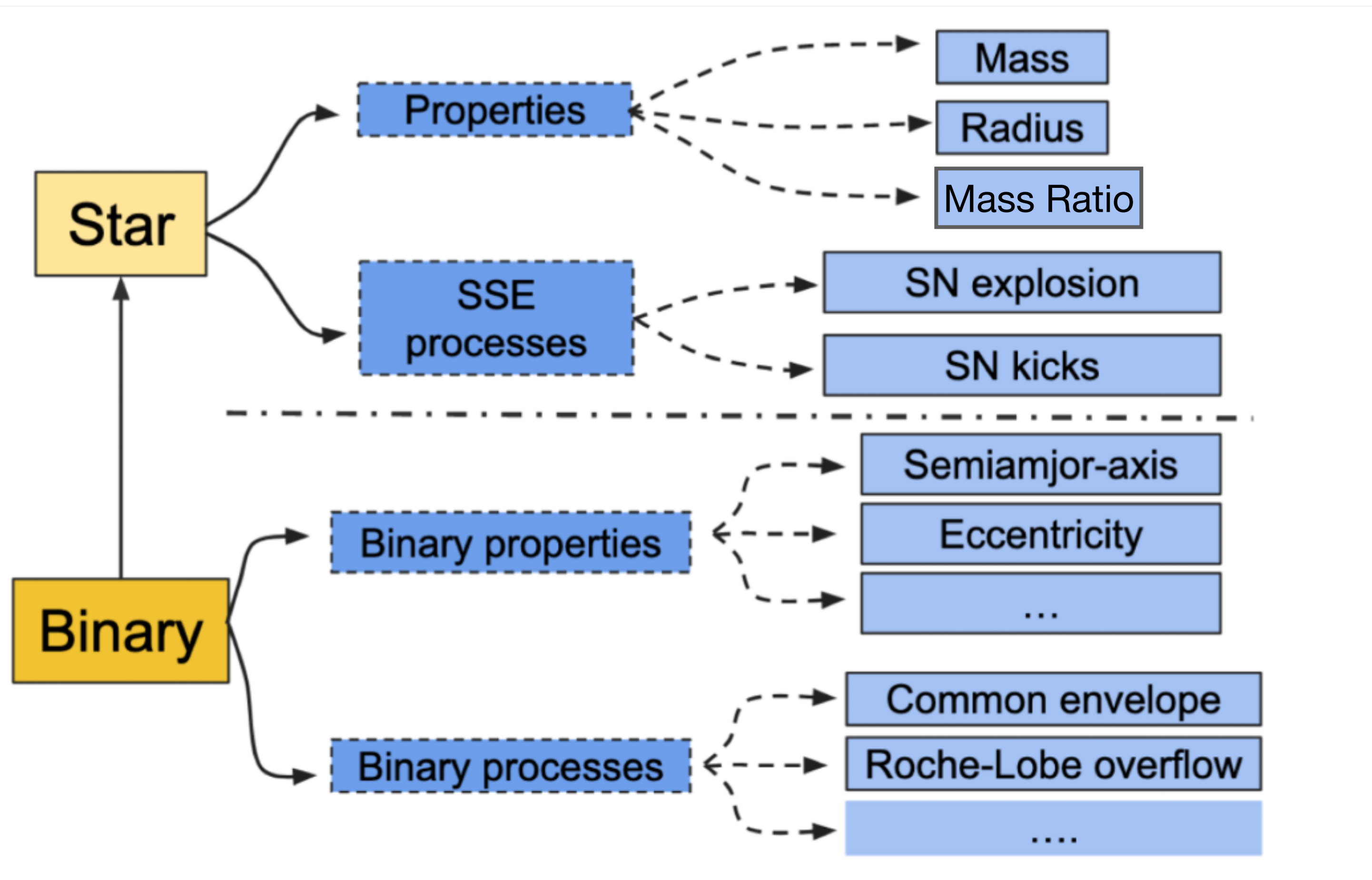


# Population synthesis

- SEVN: population-synthesis code with **stellar tracks**
- [Costa et al. 2023](#) generated a new set of Pop. III stellar tracks
- We evolve a large set of Pop III. binary stars

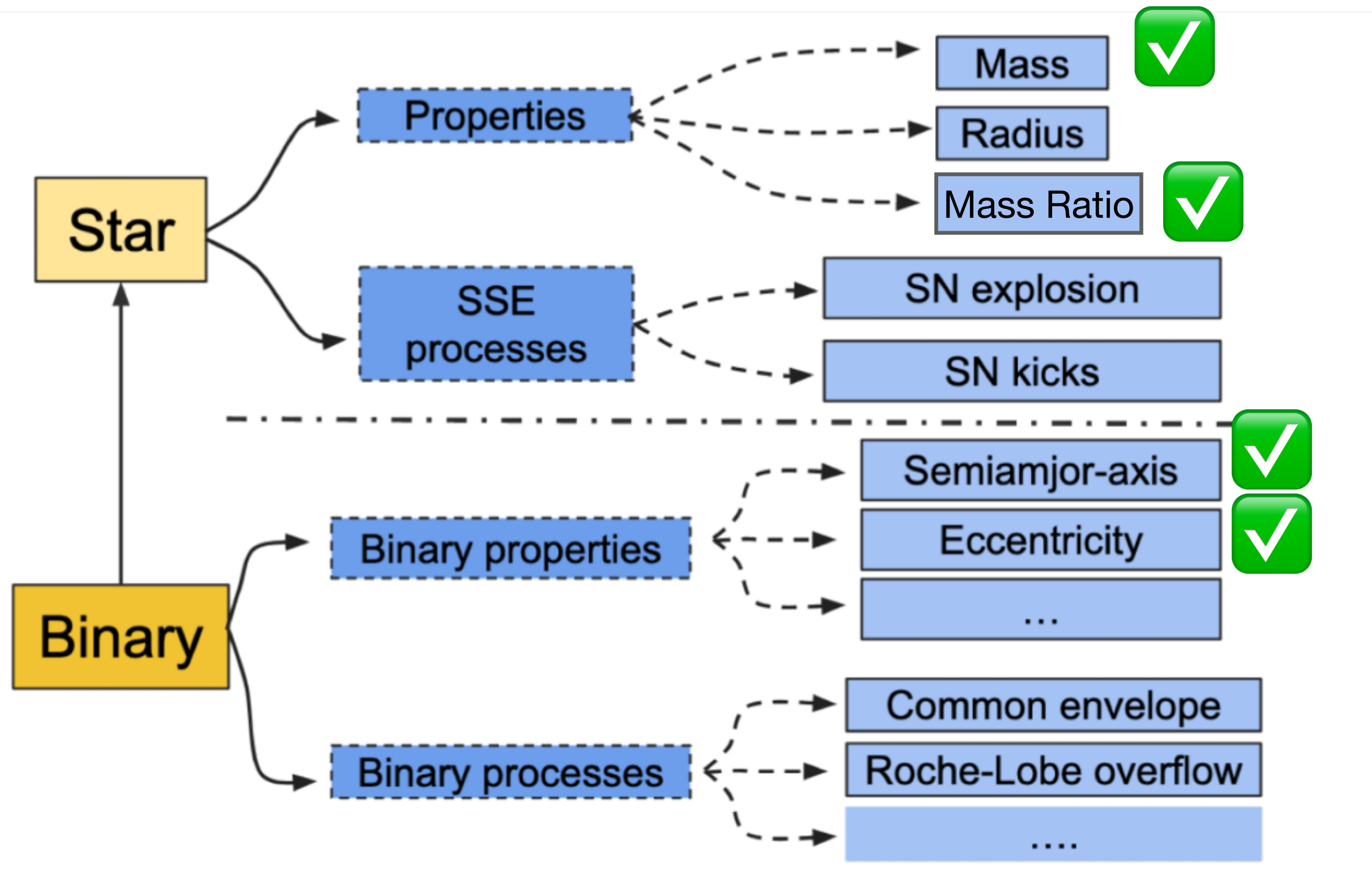


# Population synthesis



Credits: [lorio et al. 2023](#)

# Population synthesis

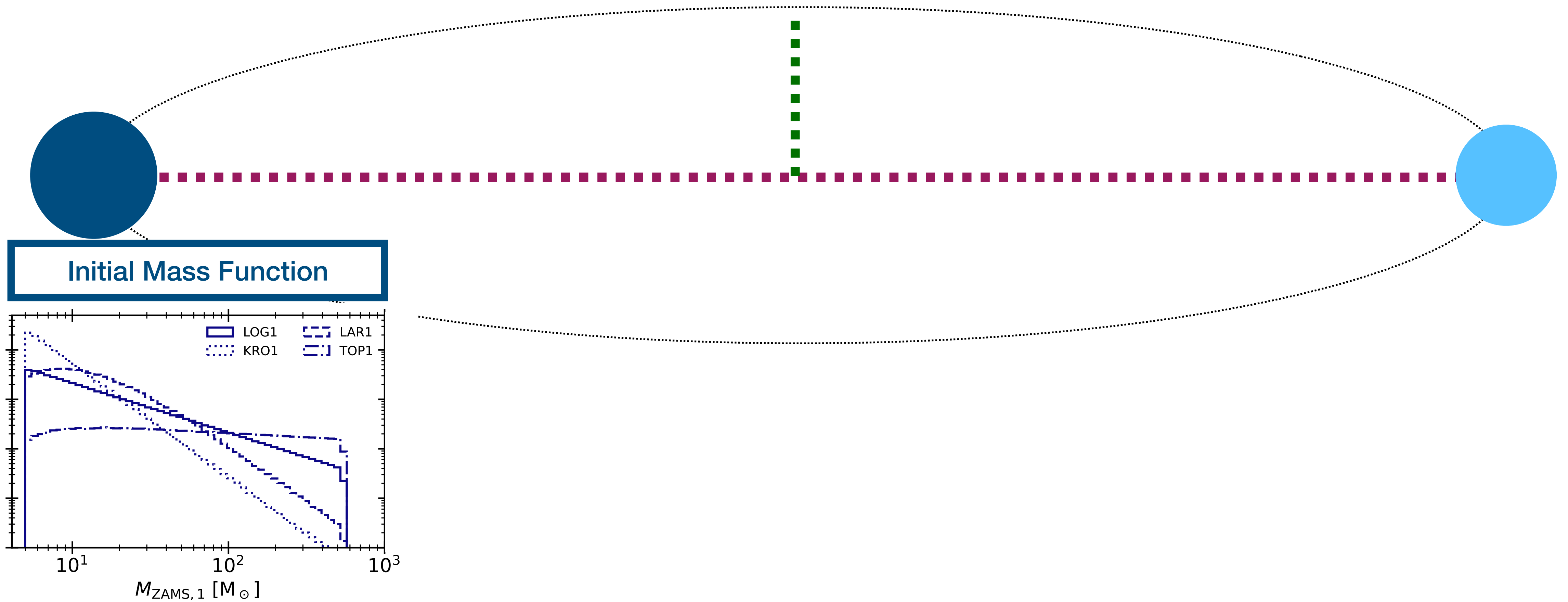


Exploring initial conditions

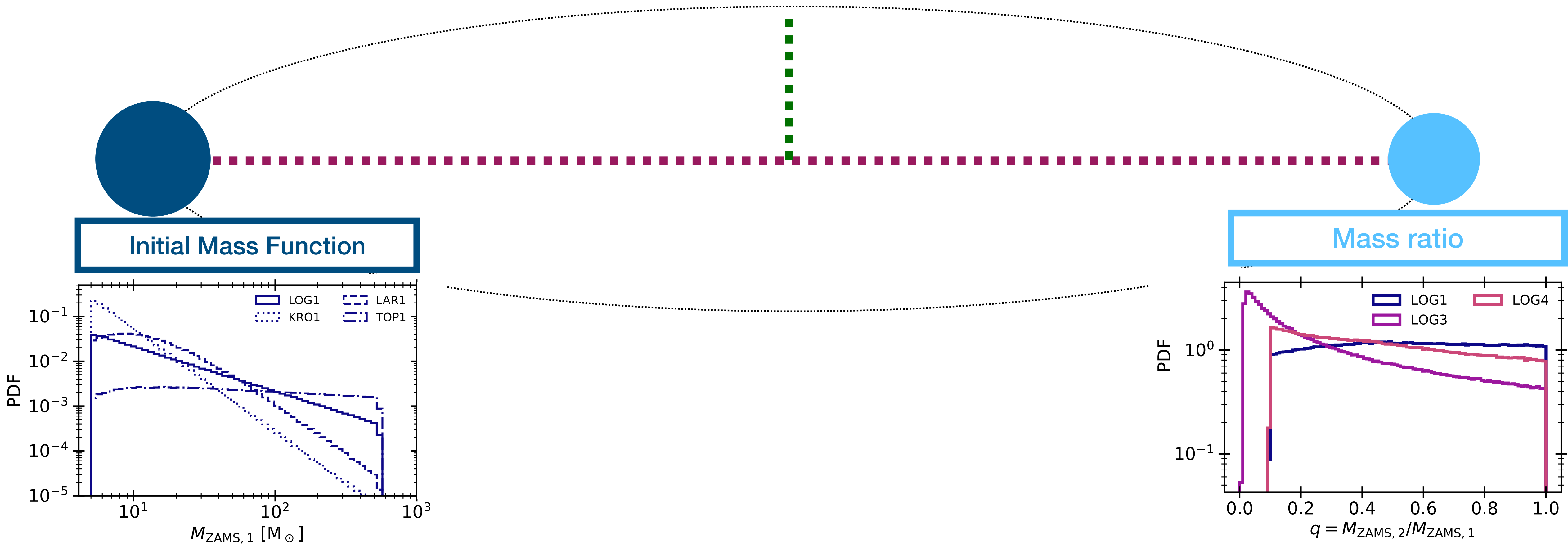
Credits: [lorio et al. 2023](#)



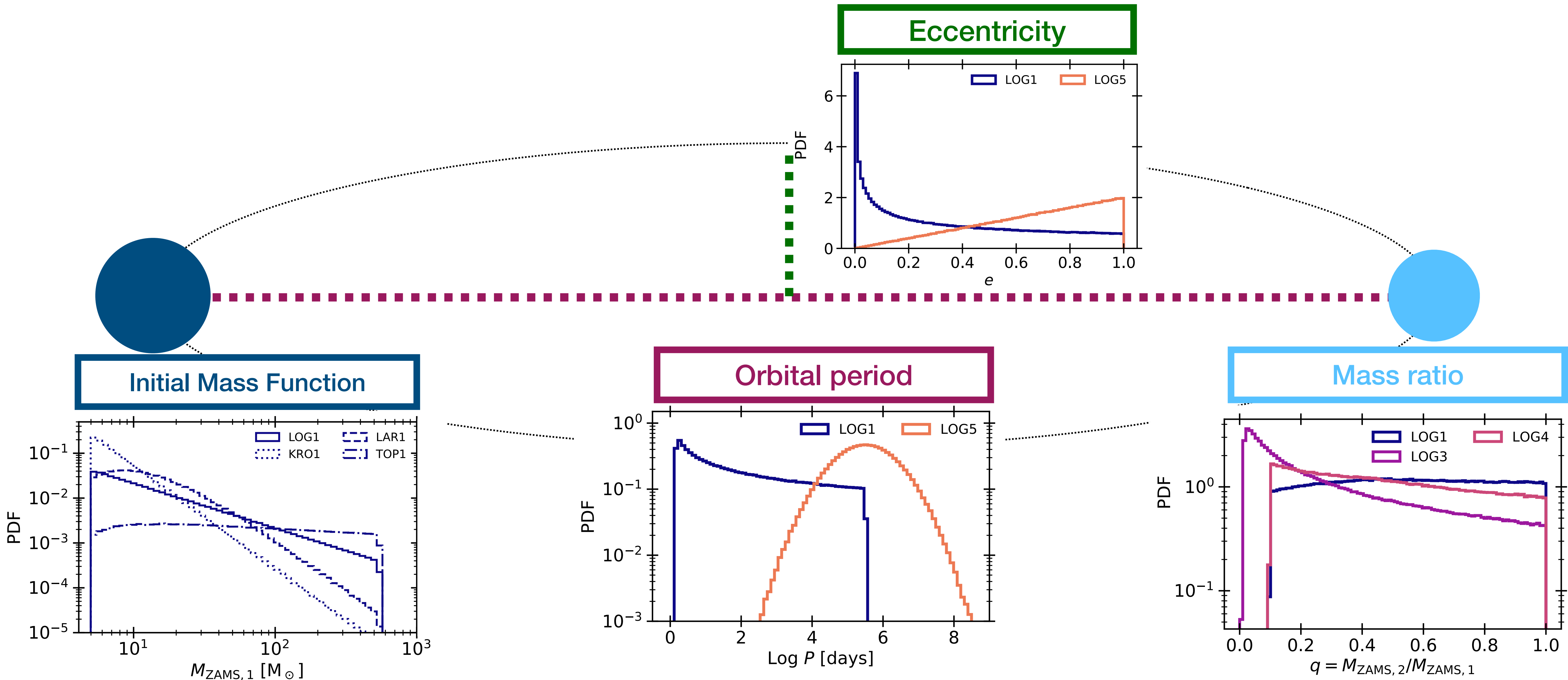
# Initial conditions



# Initial conditions



# Initial conditions





$$\mathcal{R}(z) = \int_{z_{\max}}^z \left[ \int_{Z_{\min}}^{Z_{\max}} \text{SFRD}(z', Z) \mathcal{F}(z', z, Z) dZ \right] \frac{dt(z')}{dz'} dz'$$

Output from SEVN

**Catalogs of merging BBHs:**  
primary mass, secondary mass,  
**delay time**

$$z' \rightarrow z$$

# COSMORATE

$$\mathcal{R}(z) = \int_{z_{\max}}^z \left[ \int_{Z_{\min}}^{Z_{\max}} \text{SFRD}(z', Z) \mathcal{F}(z', z, Z) dZ \right] \frac{dt(z')}{dz'} dz'$$

$$\text{SFRD}(z, Z) = \psi(z) p(Z|z)$$

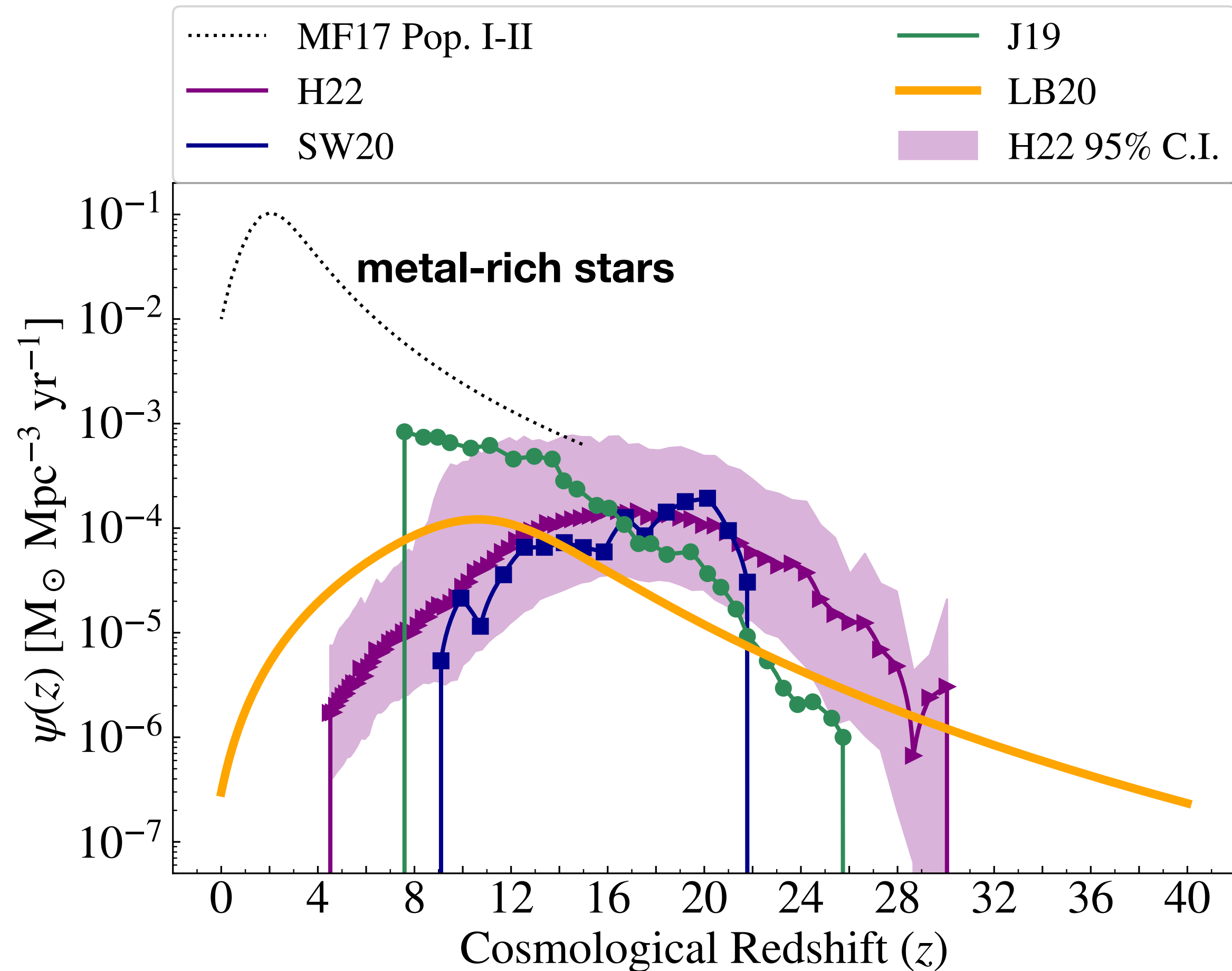
Evaluated from SEVN catalogs

$$Z = 10^{-11}$$

$$\mathcal{R}(z) = \int_{z_{\max}}^z \left[ \int_{Z_{\min}}^{Z_{\max}} \text{SFRD}(z', Z) \mathcal{F}(z', z, Z) dZ \right] \frac{dt(z')}{dz'} dz'$$

$$\text{SFRD}(z, Z) = \psi(z) p(Z|z)$$

Evaluated from SEVN catalogs



### 4 different Pop. III SFRDs:

**H22** - [Hartwig et al. 2022](#)

**J19** - [Jaacks et al. 2019](#)

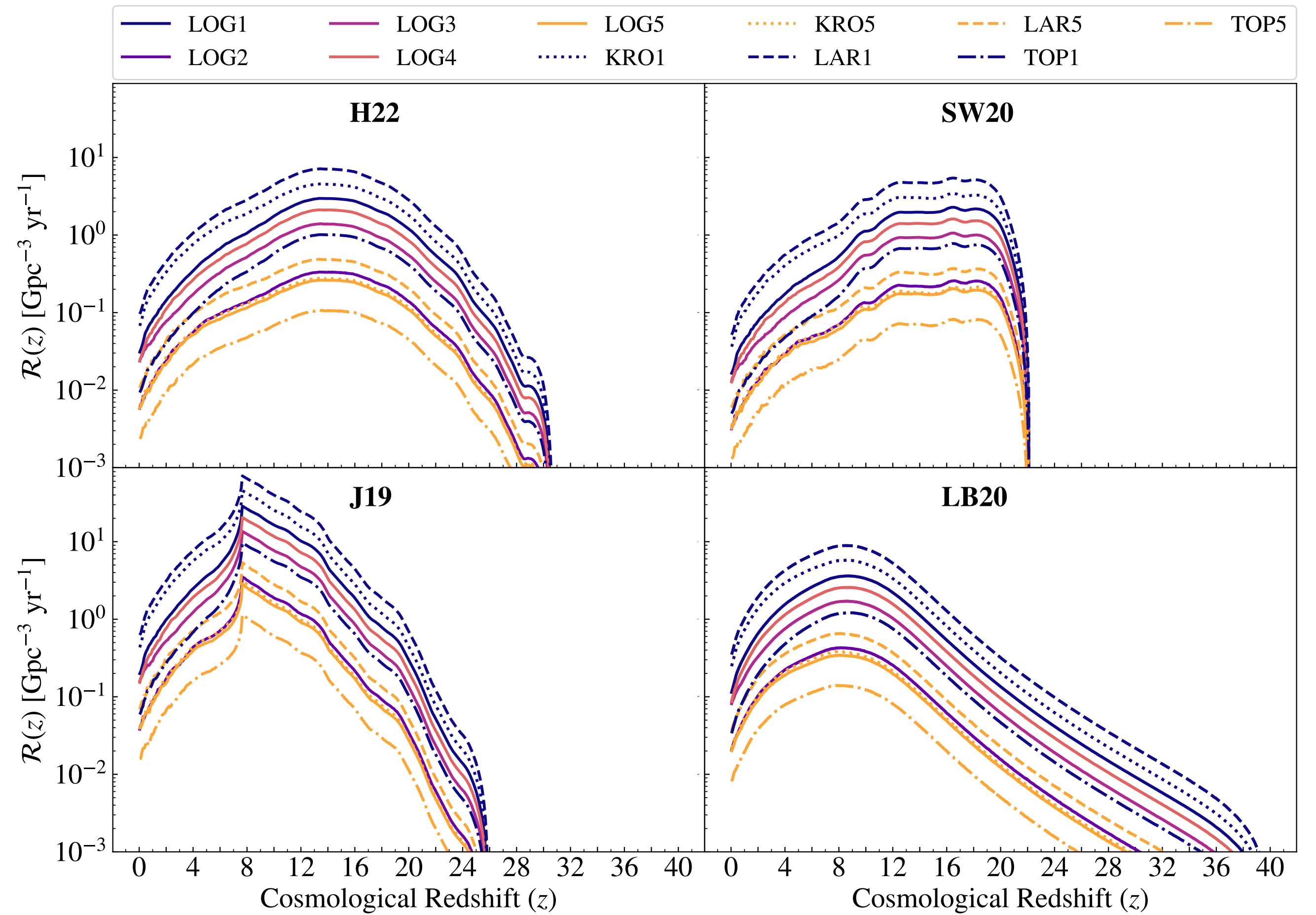
**LB20** - [Liu & Bromm 2020](#)

**SW20** - [Skinner & Wise 2020](#)

different assumptions on baryonic physics  
+ cosmic variance



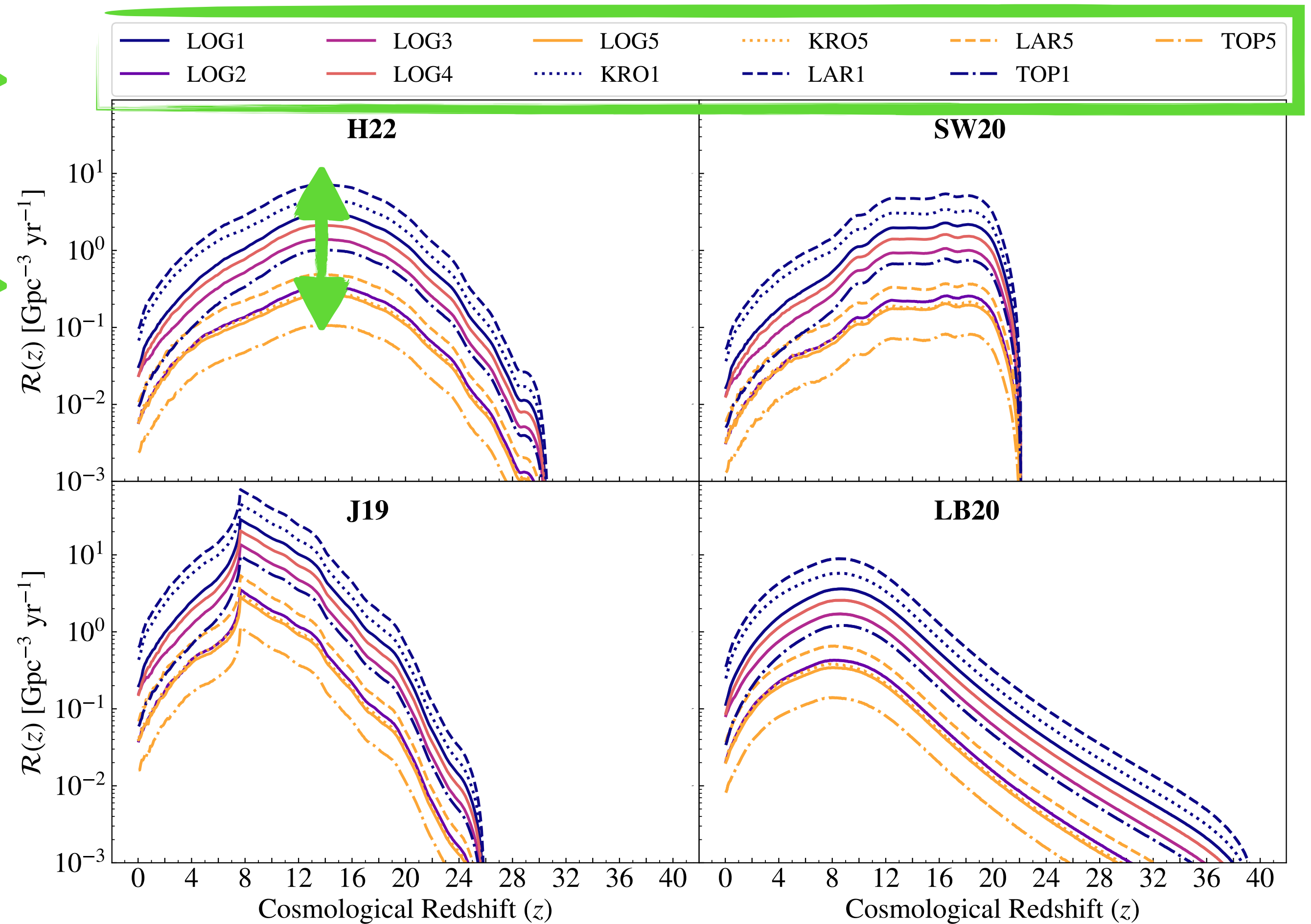
# Merger rate density



# Merger rate density

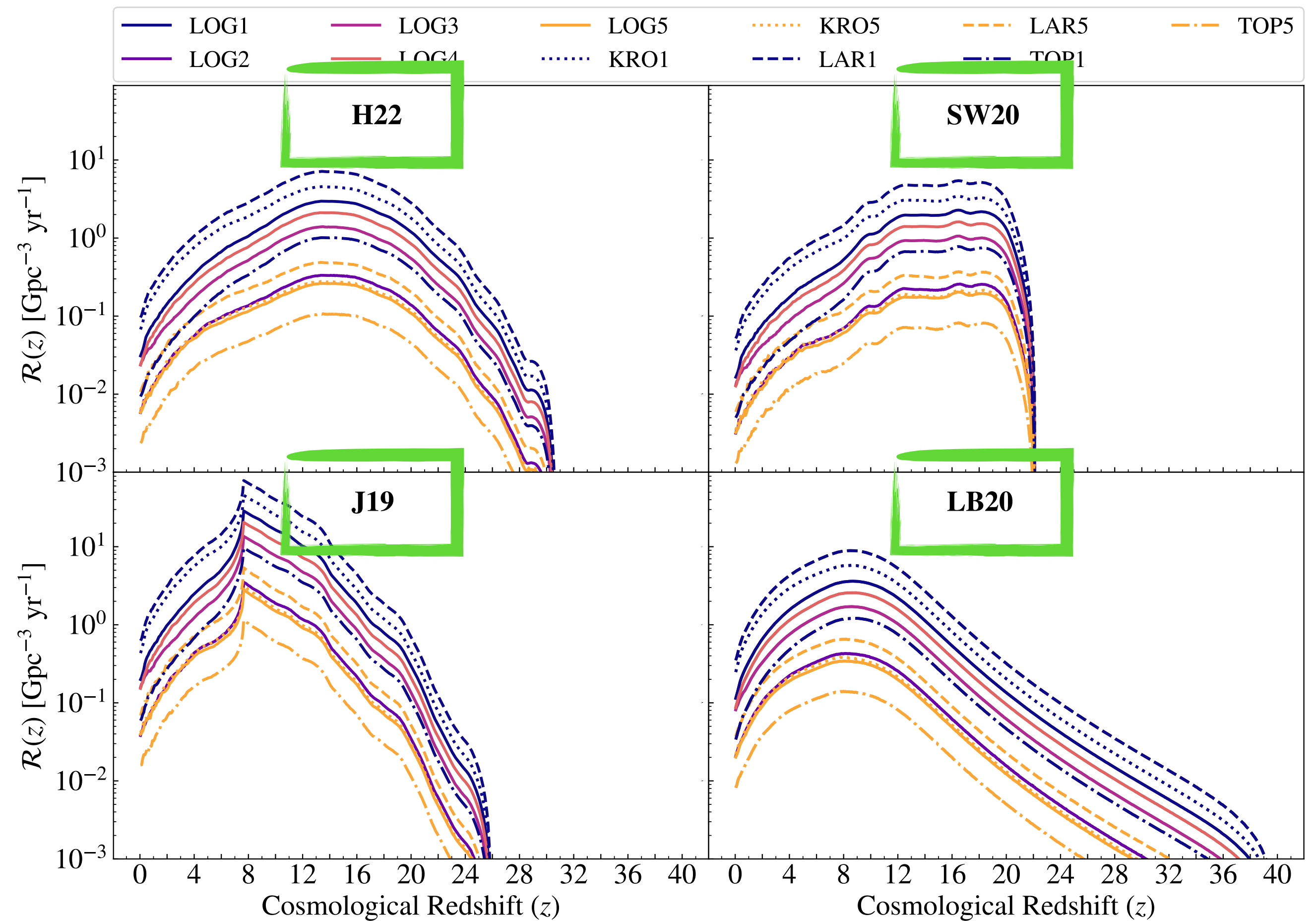
initial conditions  $\rightarrow$

impact  $\sim 2$  orders of magnitude  $\rightarrow$



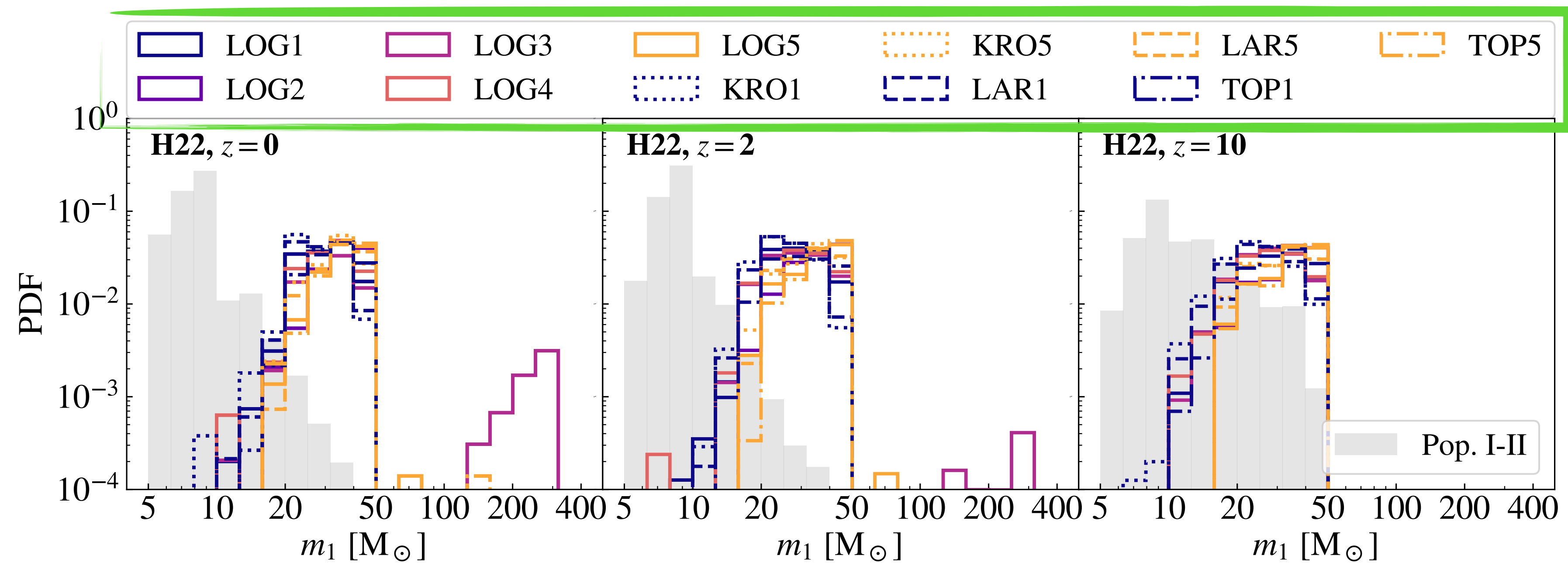
# Merger rate density

star formation history impacts shape and normalisation



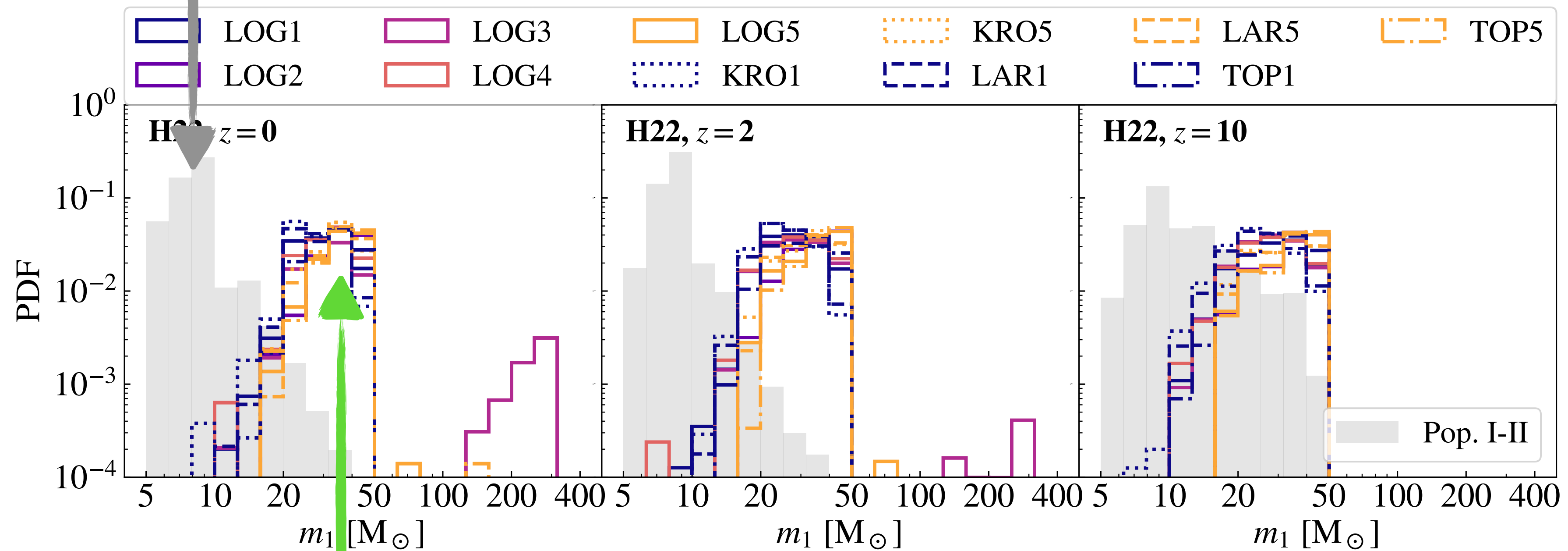


# Primary mass



# Primary mass

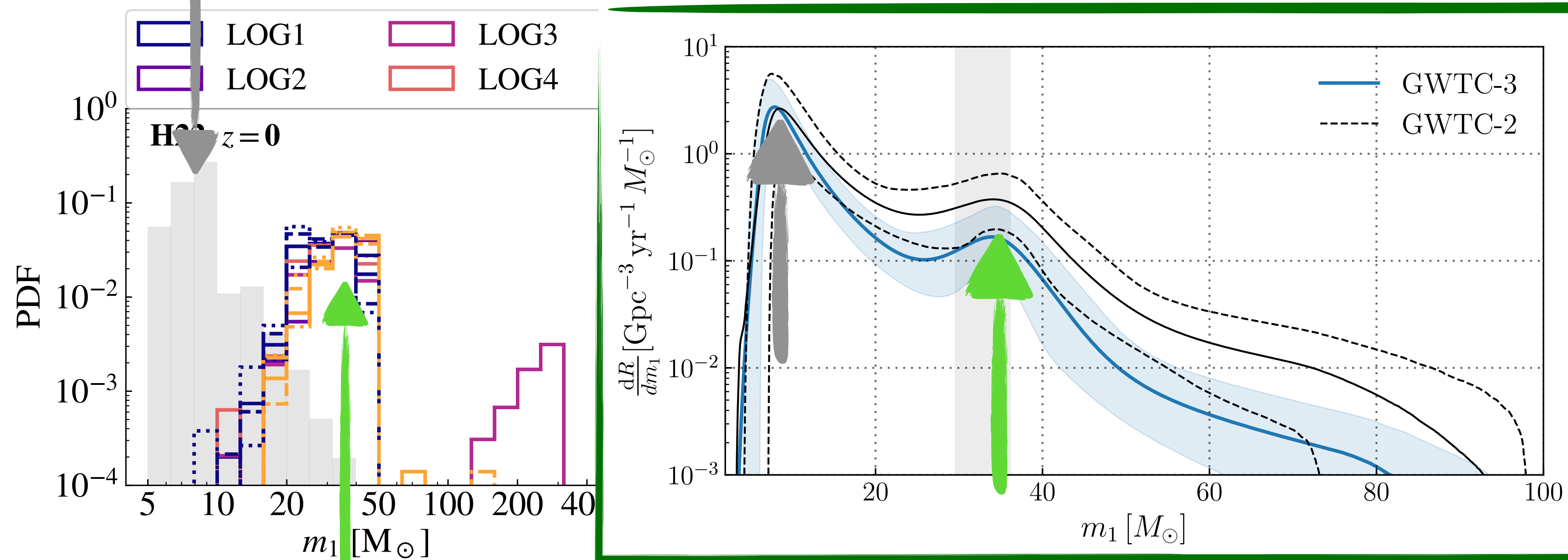
At  $z = 0$ , Pop. I-II BBHs show a main peak at 8 – 10  $M_{\odot}$



Pop. III BBHs show a main peak at 30 – 35  $M_{\odot}$

# Primary mass

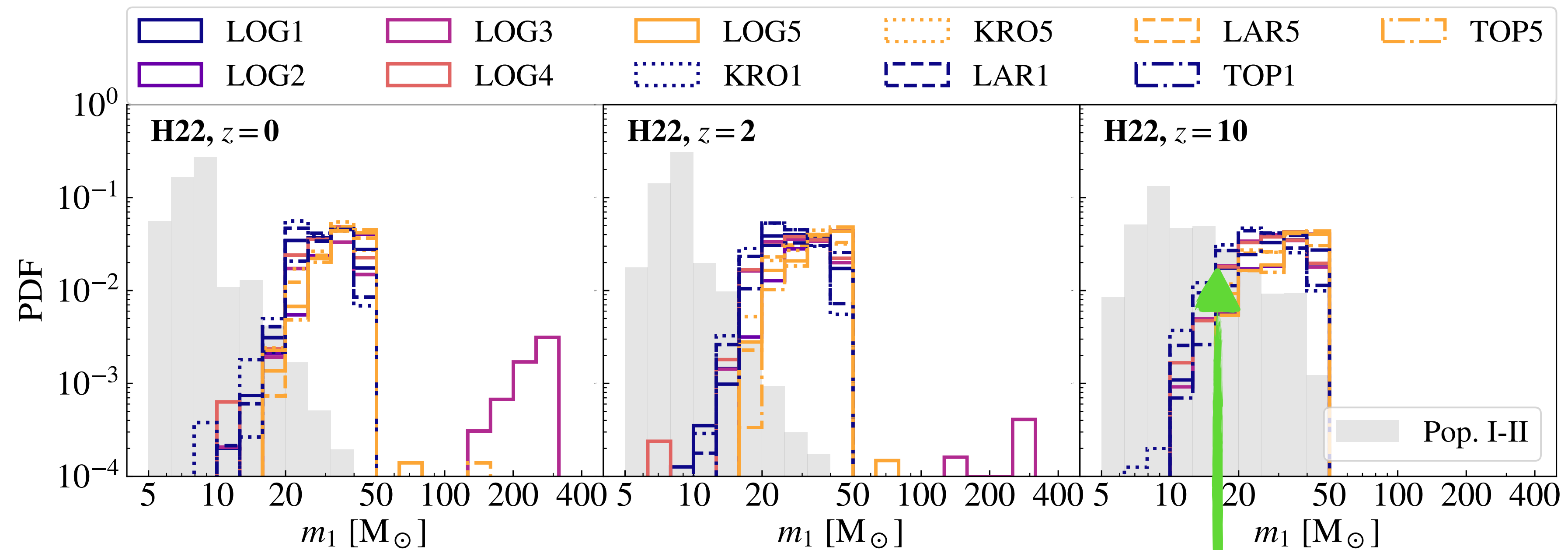
At  $z = 0$ , Pop. I-II BBHs show a main peak at 8 – 10  $M_{\odot}$



Pop. III BBHs show a main peak at 30 – 35  $M_{\odot}$



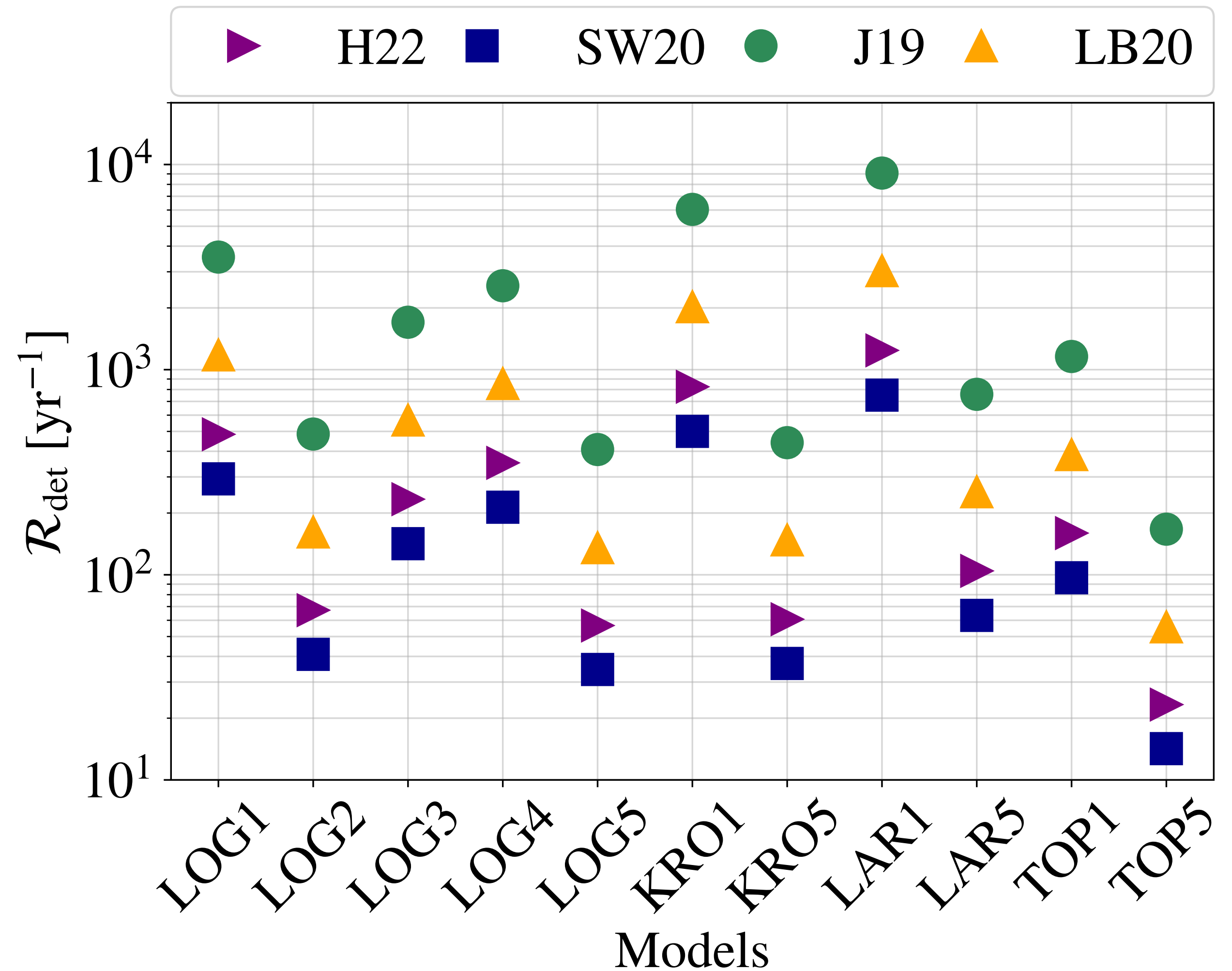
# Primary mass



At high redshift, overlap increases

# Detection rate

- Einstein Telescope will detect  $10 - 10^4$  Pop. III BBH mergers per year
- We expect between 23% and 73% of detections to occur at redshift  $z > 8$



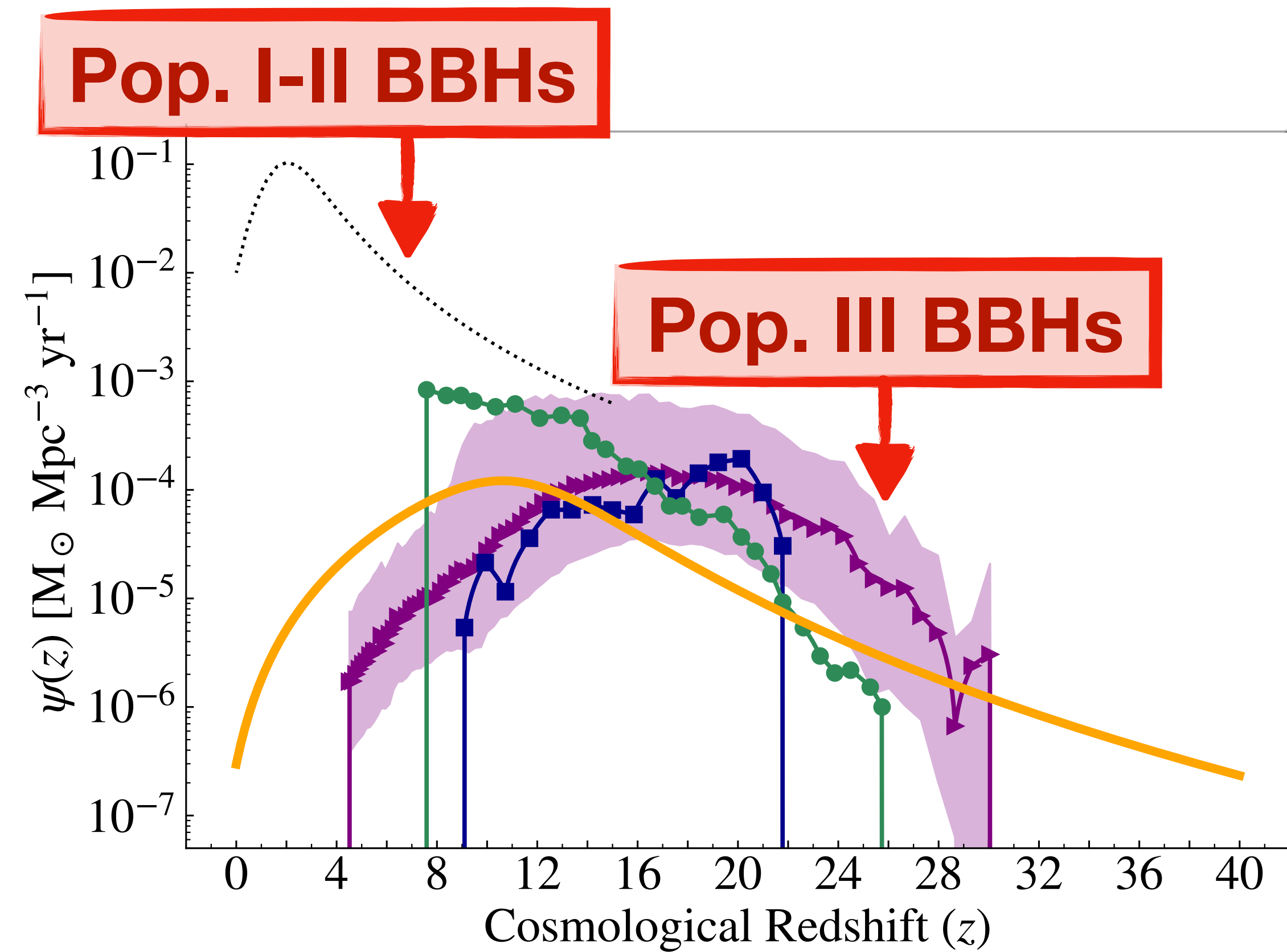
# Can we distinguish Pop. III BBHs?

- Einstein Telescope will detect BBH mergers up to  $z \sim 100$
- high-redshift sources with low-SNR and poor estimate of  $d_L$
- inferring the origin of individual GW detections will not be granted

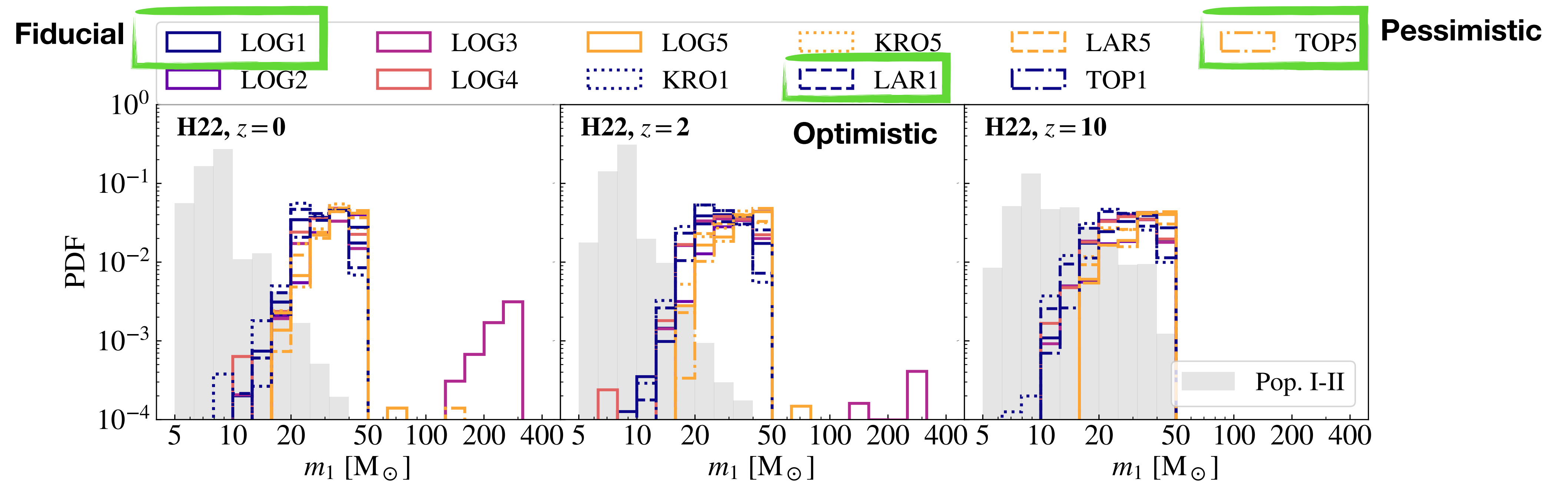
Ref. [Maggiore et al. 2020](#), [Ng et al. 2021](#), [2022](#), [Branchesi et al. 2023](#), [Mancarella et al. 2023](#)



# Goal: classification

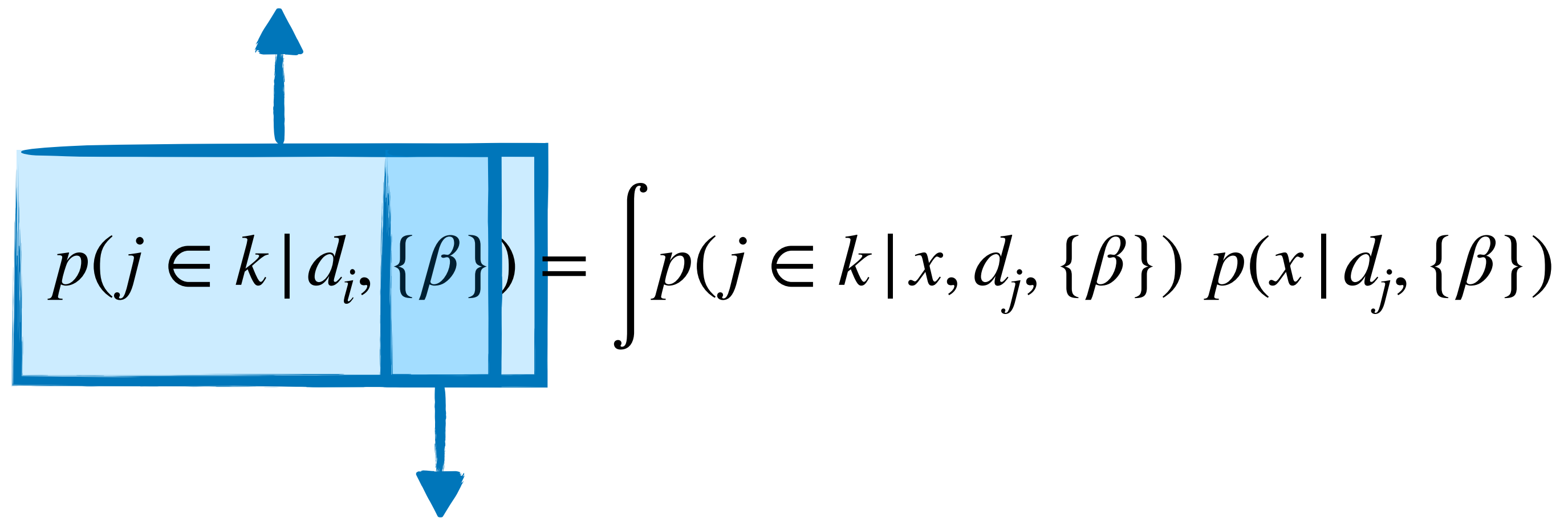


# Simulation-based classification



# Classification

This is the probability that the event  $j$  is a Pop. III BBH


$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$

**Mixing fraction**

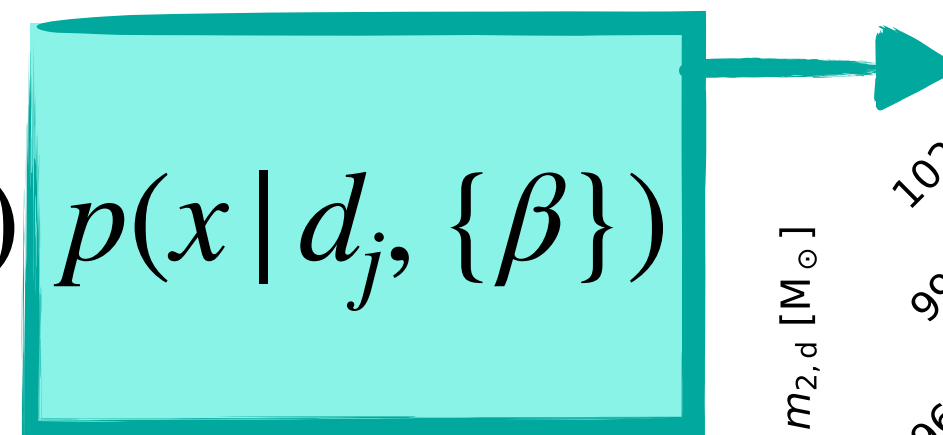
**Fiducial ~4%**  $\beta_{III} \propto N_{\text{Pop. III}} \sim 400$

$\beta_{I-II} \propto N_{\text{Pop. I-II}} \sim 10^4$

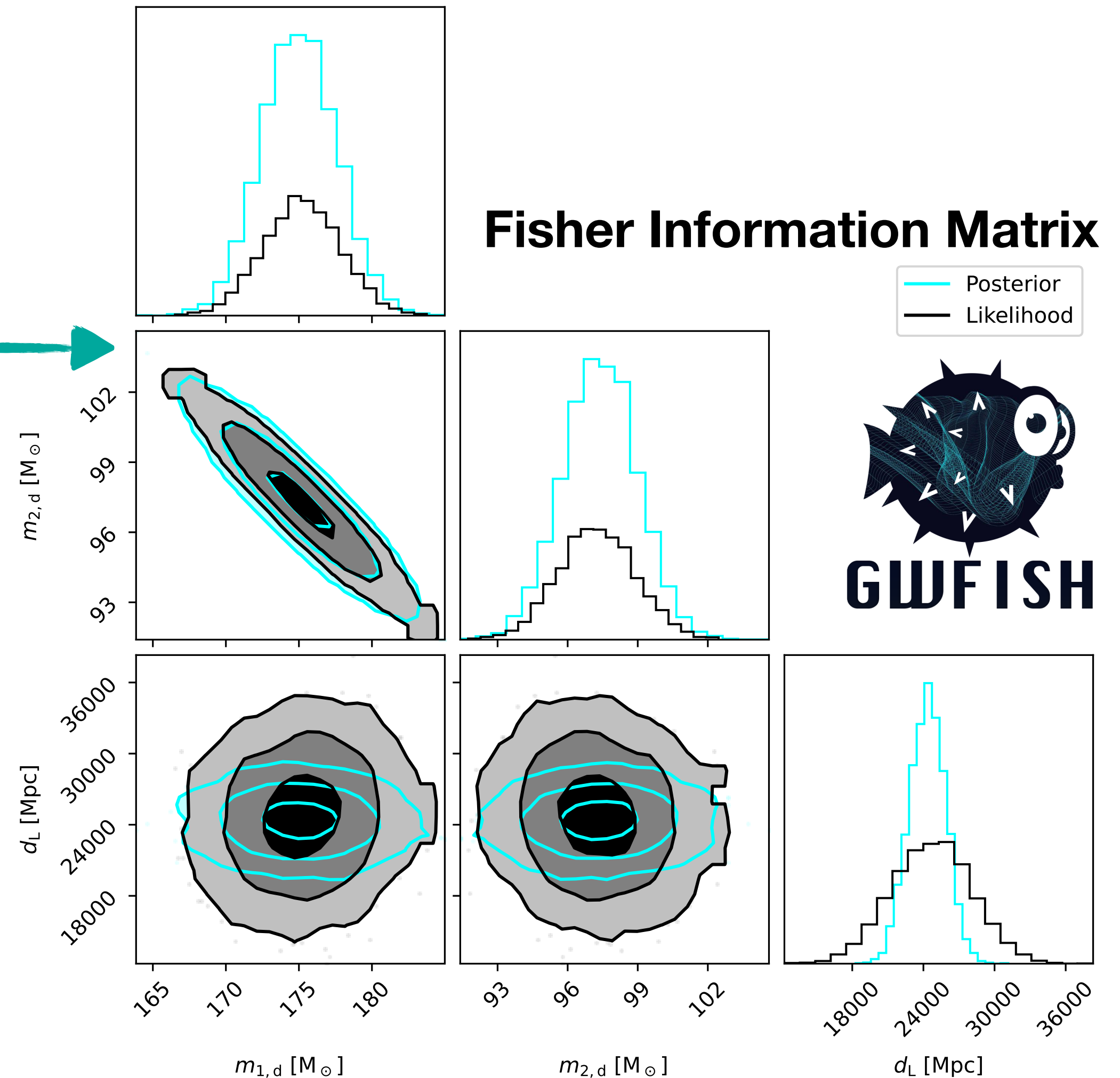


# Classification

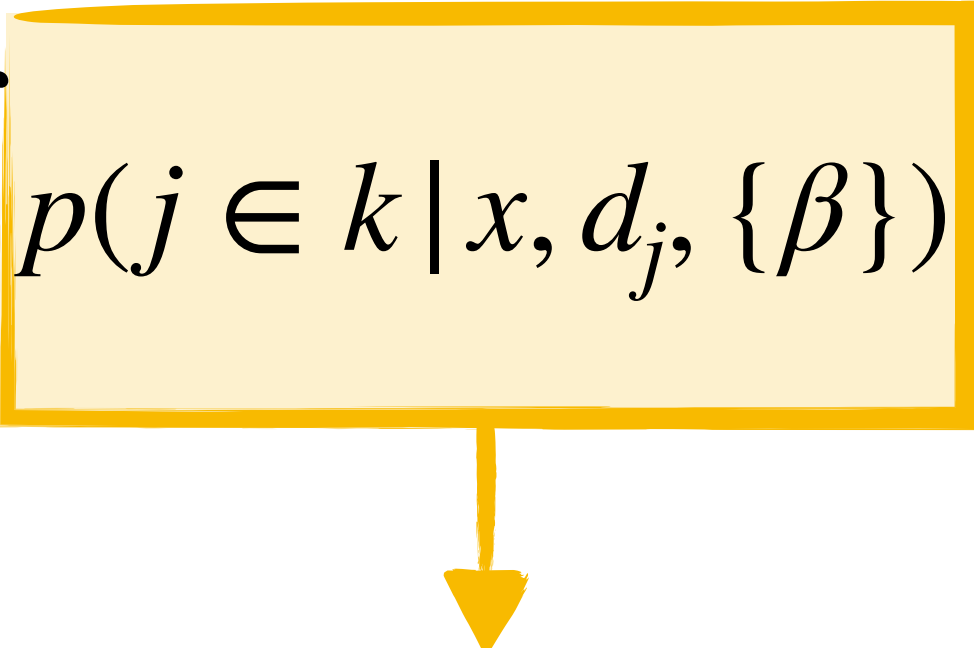
$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$



This is the posterior of waveform parameters →  
parameter estimation **performance of ET**




# Classification

$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$


This is the probability that links waveform parameters to Pop. III BBHs

➡ easy to consider a fix threshold

$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$

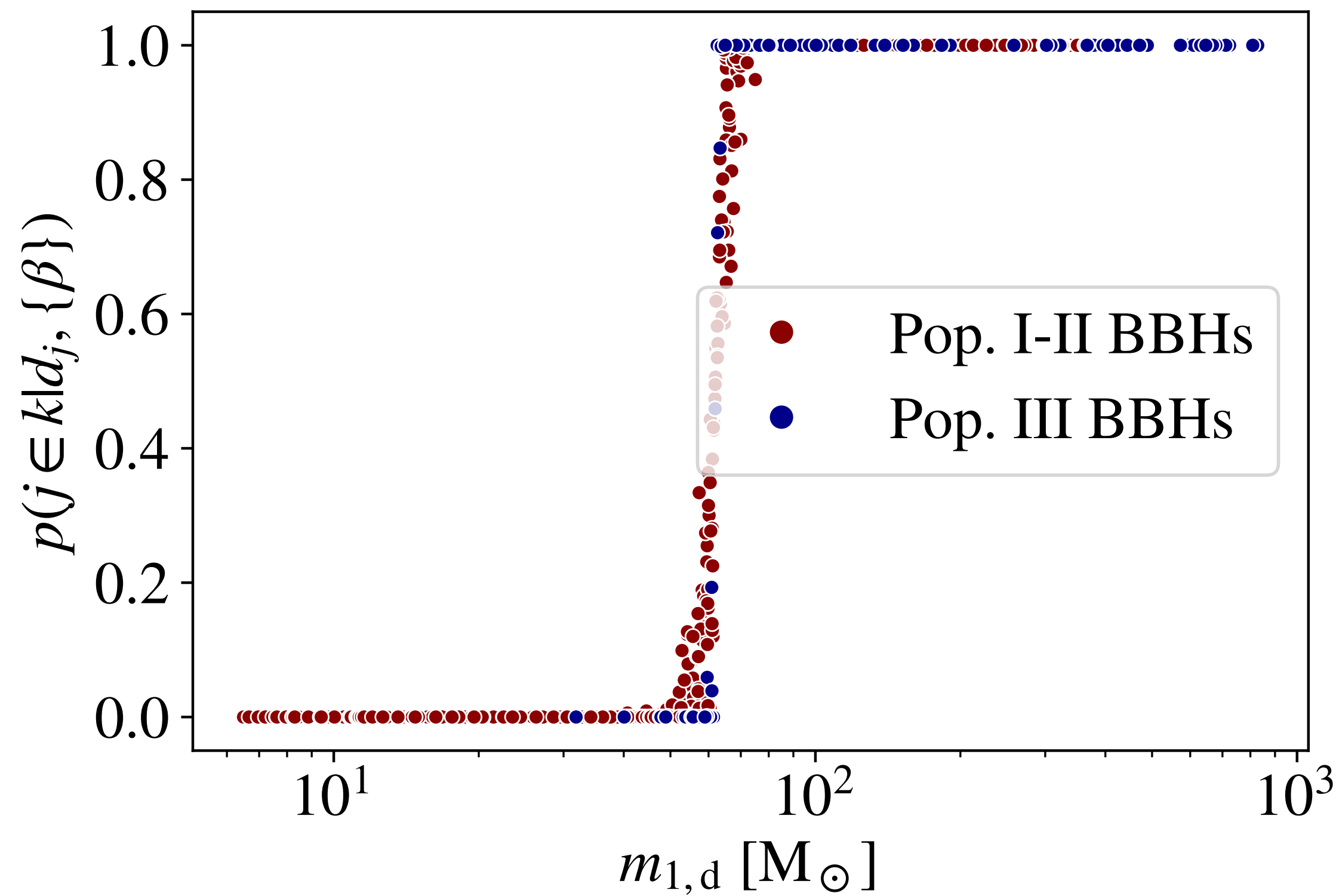


$$= 1 \quad \text{if } m_{1,d} \gtrsim 60 M_{\odot}$$
$$m_{1,d} = m_1(1 + z)$$

$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$

$$= 1 \quad \text{if} \quad m_{1,d} \gtrsim 60 M_{\odot}$$

manual classifier - fiducial



⚠ low performances: precision is  $\sim 0.16$



$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$



we can use Machine Learning

***XGBoost***

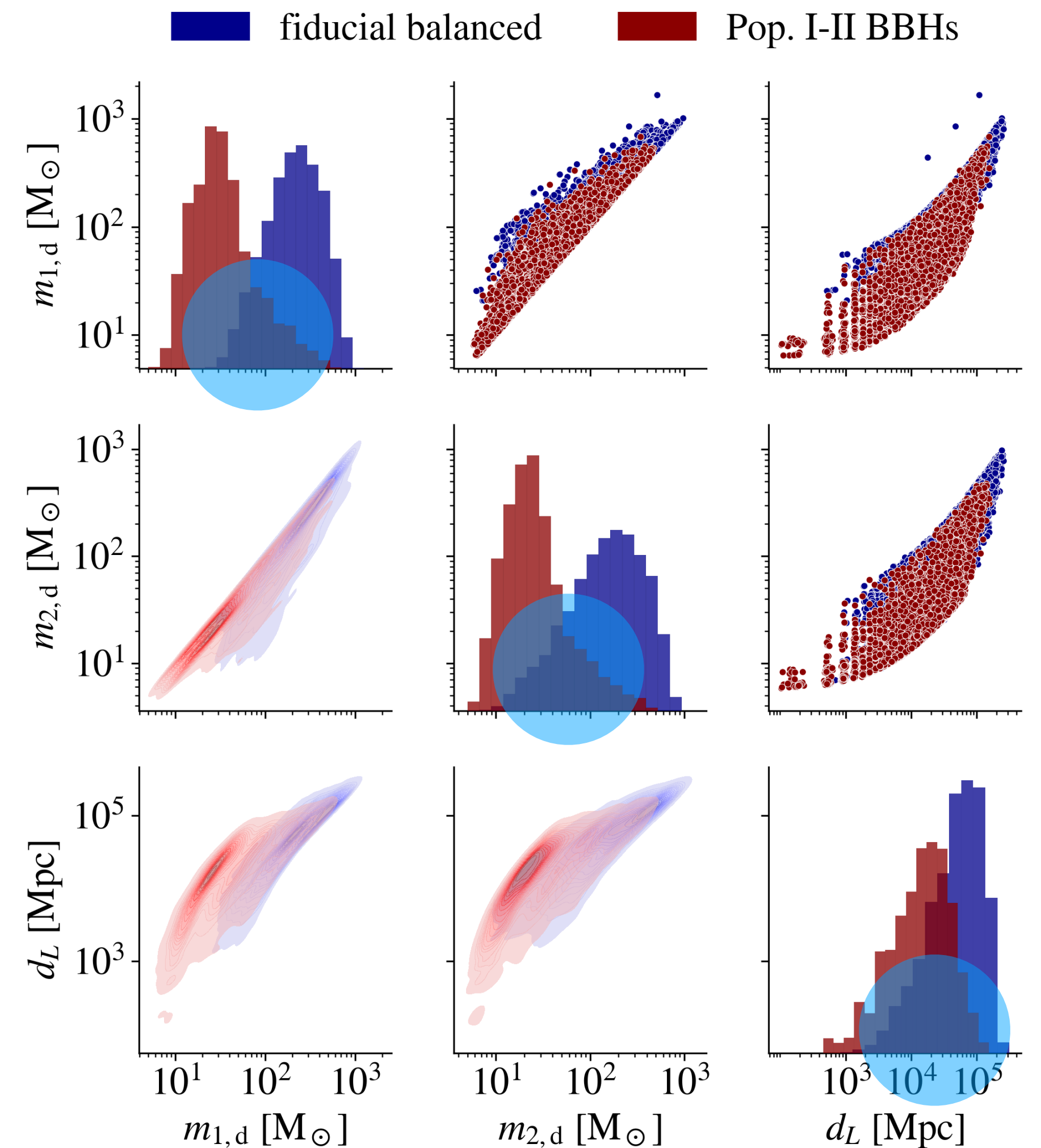
supervised ML based on decision trees

$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$

Using Machine Learning

\* trained and tested on balanced classes + re-balancing

\* instances:  $> 10^4$



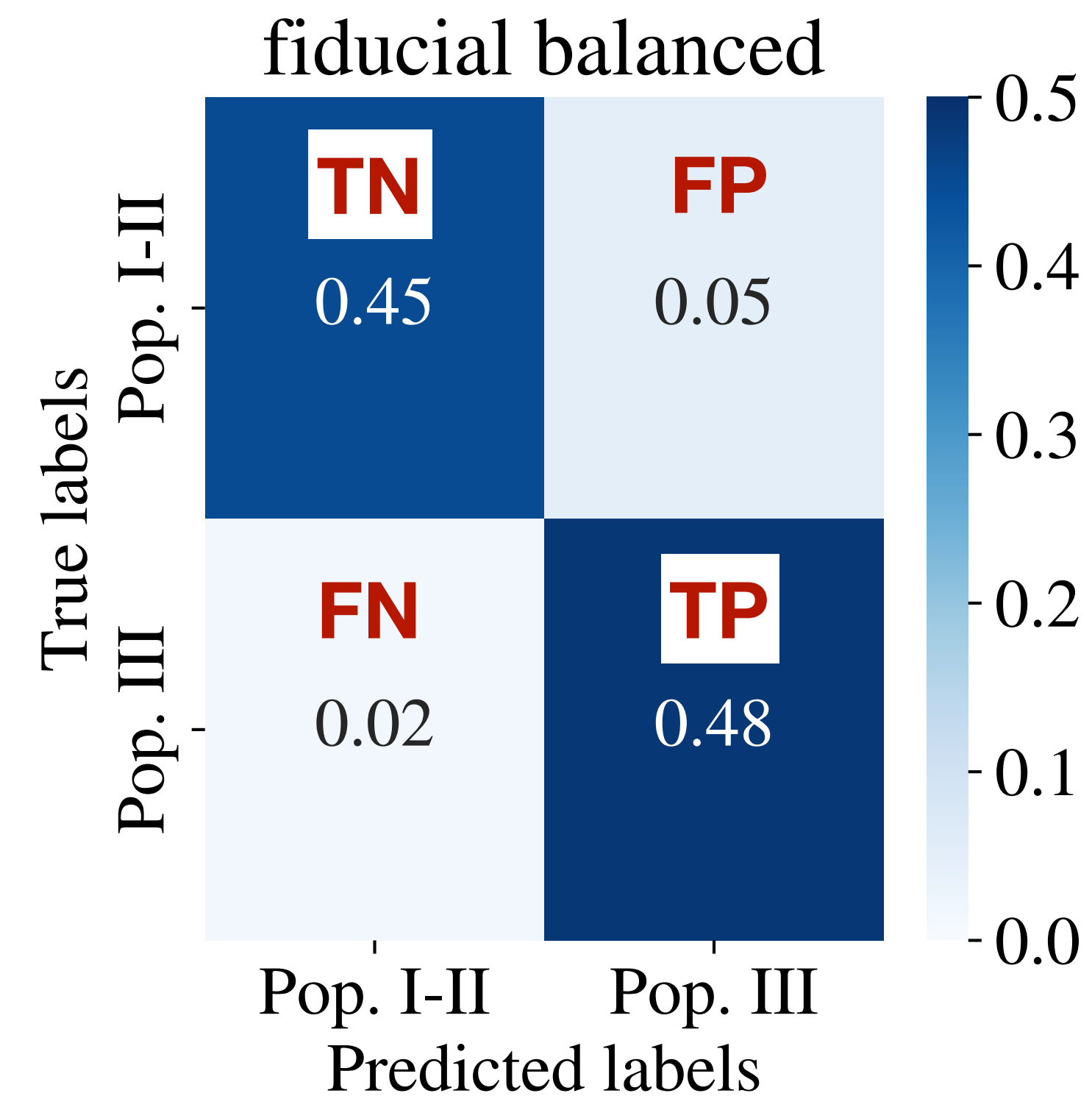
$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$



Using Machine Learning



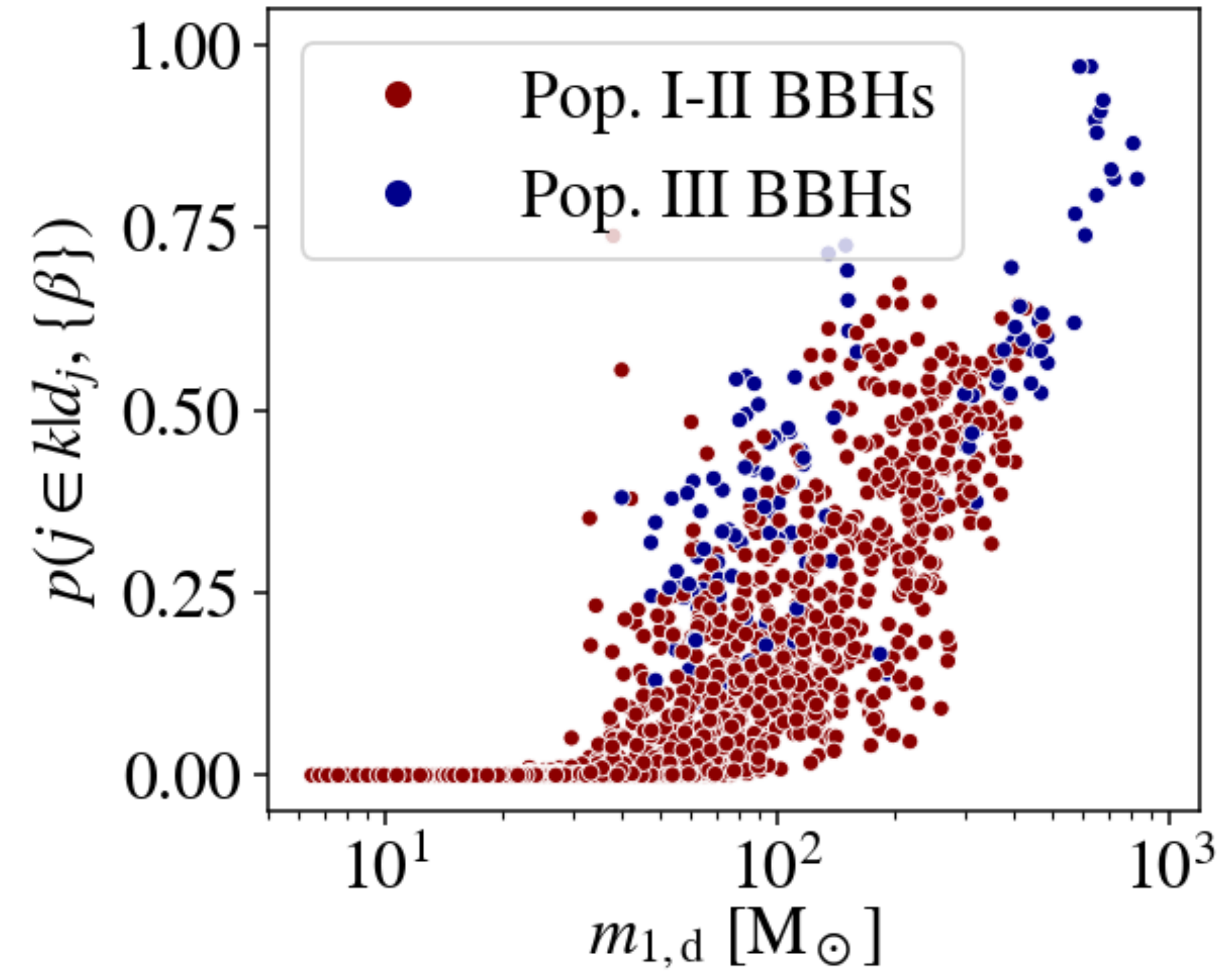
**precision = TP/(TP+FP)**  
**high precision > 0.90**



$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$

Using Machine Learning

fiducial



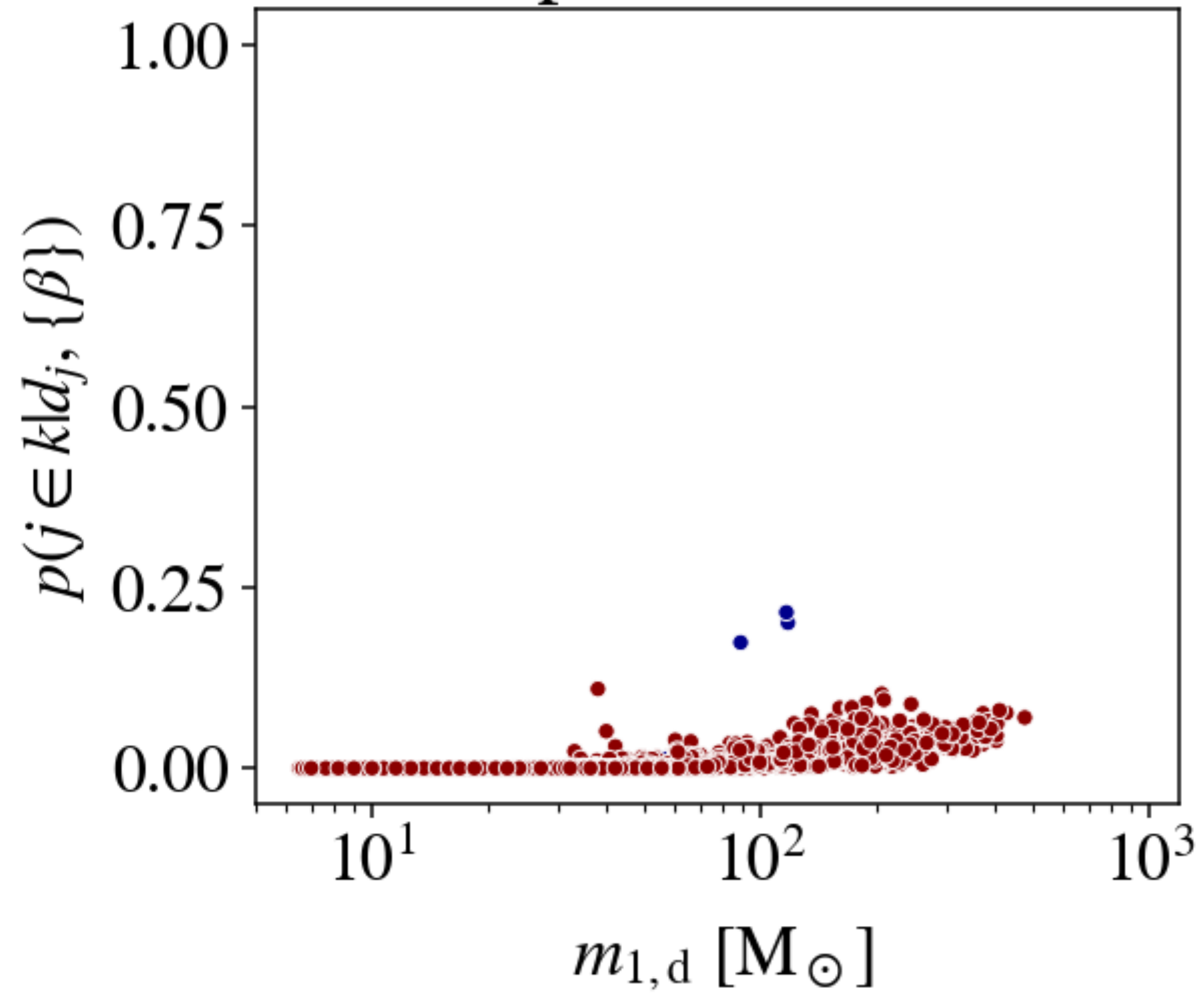
**~10% of detected sources are classified with precision > 0.90**



$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$

Using Machine Learning

pessimistic

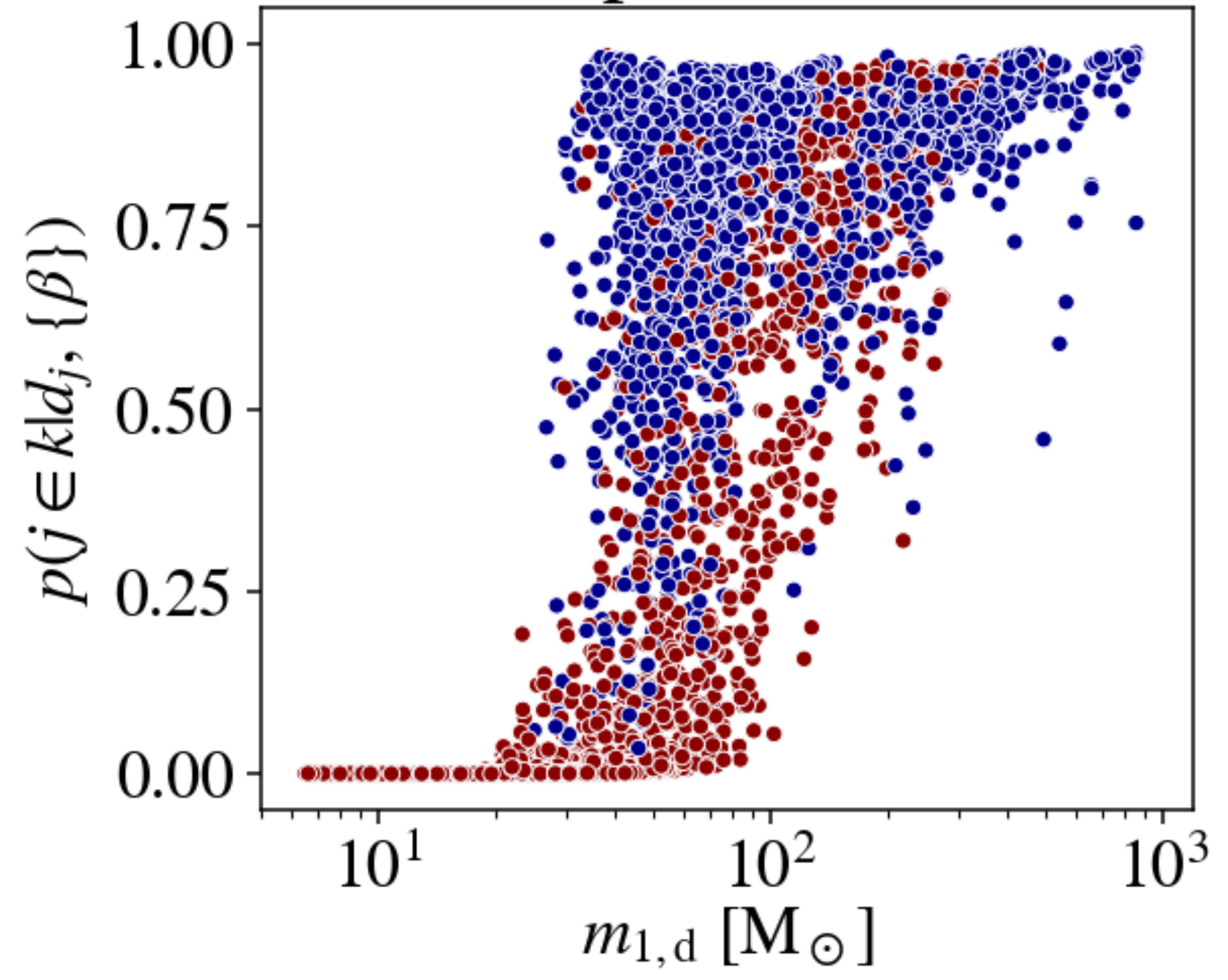


**~30% of detected sources are classified with precision > 0.90**

$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$

Using Machine Learning

optimistic



**~45% of detected sources are classified with precision > 0.90**

# Contributions

- First large parameter exploration of Pop. III BBHs
  - SFRD affects normalisation and shape of merger rate density
  - primary mass of Pop. III BHs is substantially massive
- ET will detect these sources and machine learning boosts our ability to classify them

# Backup slides



# $\alpha\lambda$ formalism for modelling the common envelope

- $\Delta E = \alpha(E_{b,f} - E_{b,i}) = \alpha \frac{Gm_{c1}m_{c2}}{2} \left( \frac{1}{a_f} - \frac{1}{a_i} \right)$  This is the orbital energy before and after the common envelope phase
- $E_{\text{env}} = \frac{G}{\lambda} \left[ \frac{m_{\text{env},1}m_1}{R_1} + \frac{m_{\text{env},2}m_2}{R_2} \right]$  This is the binding energy of the envelope
- By imposing  $\Delta E = E_{\text{env}}$ ,  $\frac{1}{a_f} = \frac{1}{\alpha\lambda} \frac{2}{m_{c1}m_{c2}} \left[ \frac{m_{\text{env},1}m_1}{R_1} + \frac{m_{\text{env},2}m_2}{R_2} \right] + \frac{1}{a_i}$
- If  $\alpha$  is larger,  $a_f$  is larger, following  $a_f \sim \frac{\alpha}{1 + \alpha}$ . Therefore larger  $\alpha$  gets wider binaries
- Where  $\lambda$  is the parameter which measures the concentration of the envelope (the smaller  $\lambda$  is, the more concentrated is the envelope).
- The  $\alpha\lambda$  formalism is a simplified prescription. When  $\alpha > 1$ , we account for other sources of energy that make the envelope less bind, for instance recombination energy. Recent works (e.g. [Fragos et al. 2019](#)) suggest that  $\alpha > 1$  is necessary to reproduce the final orbital separation obtained with hydrodynamical simulations.

# Initial conditions

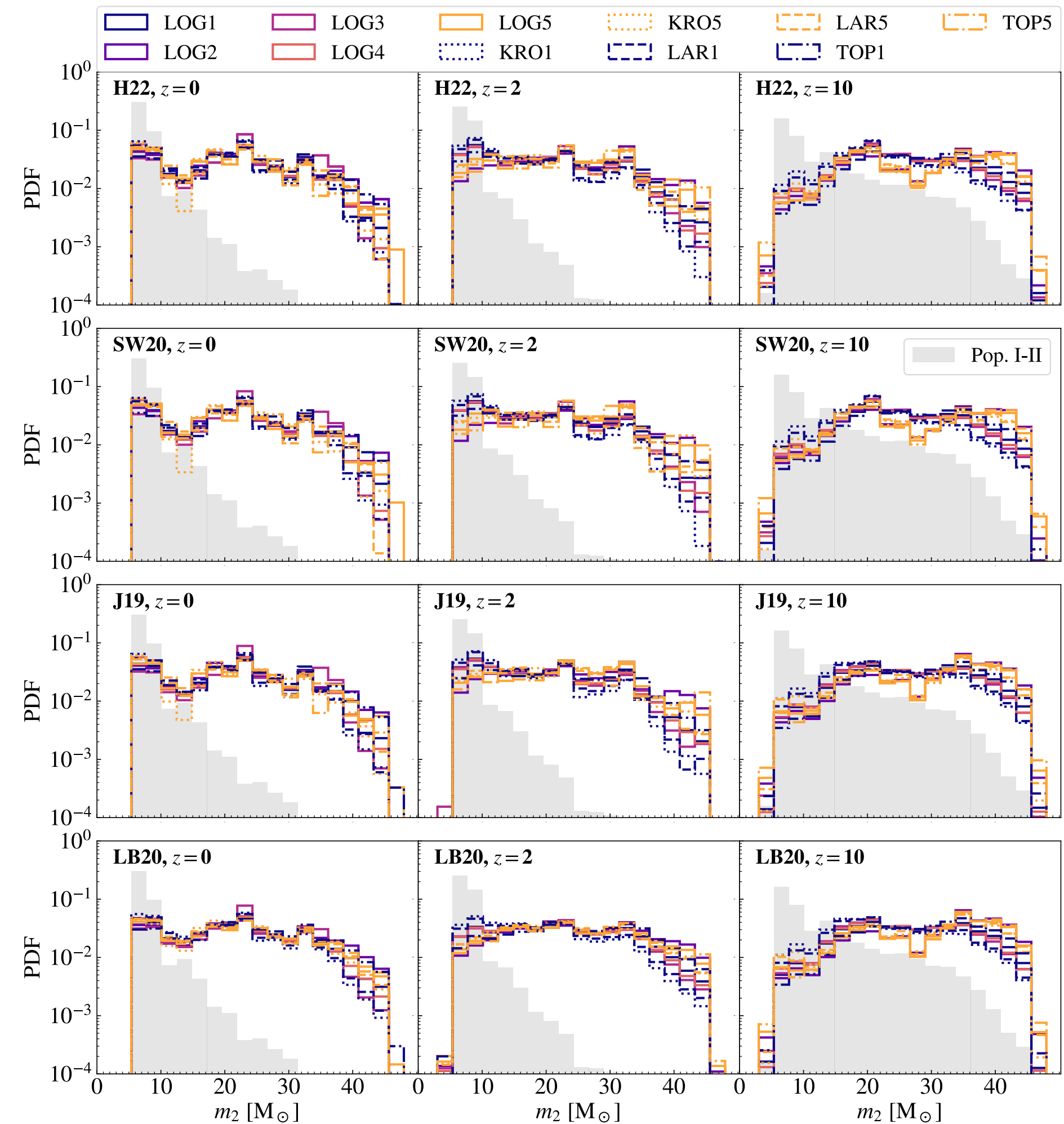
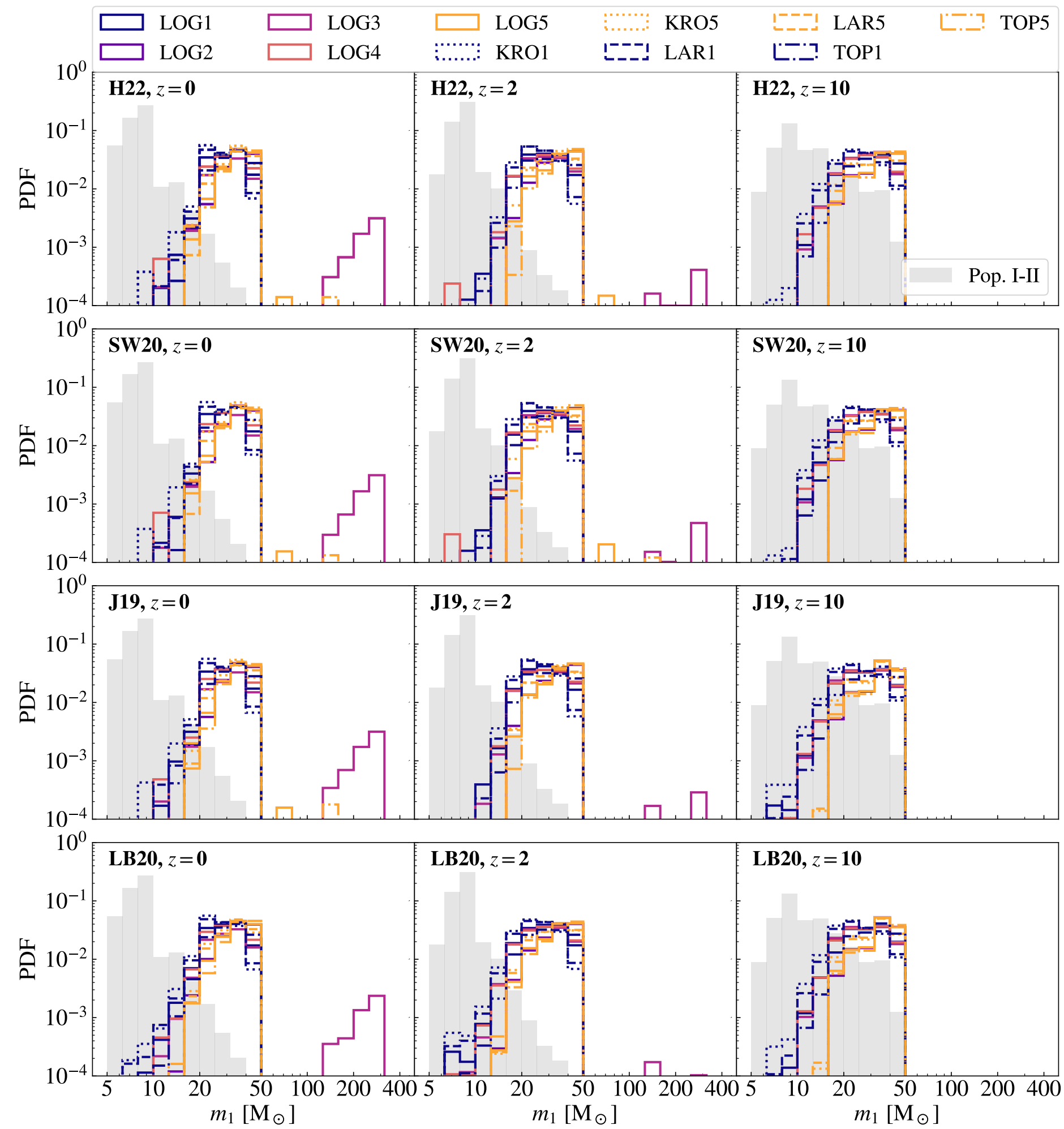
**Table 1.** Initial conditions.

Model	$M_{\text{ZAMS},1}$	$M_{\text{ZAMS}}$	$q$	$P$	$e$
LOG1	Flat in log	–	S12	S12	S12
LOG2	Flat in log	–	S12	SB13	Thermal
LOG3	–	Flat in log	Sorted	S12	S12
LOG4	Flat in log	–	SB13	S12	Thermal
LOG5	Flat in log	–	SB13	SB13	Thermal
KRO1	K01	–	S12	S12	S12
KRO5	K01	–	SB13	SB13	Thermal
LAR1	L98	–	S12	S12	S12
LAR5	L98	–	SB13	SB13	Thermal
TOP1	Top heavy	–	S12	S12	S12
TOP5	Top heavy	–	SB13	SB13	Thermal

Column 1 reports the model name. Column 2 describes how we generate the ZAMS mass of the primary star (i.e., the most massive of the two members of the binary system). Column 3 describes how we generate the ZAMS mass of the overall stellar population (without differentiating between primary and secondary stars). We follow this procedure only for model LOG3 (see the text for details). Columns 4, 5, and 6 specify the distributions we used to generate the mass ratios  $q$ , the orbital periods  $P$  and the orbital eccentricity  $e$ . See Section 2.2 for a detailed description of these distributions.

**Santoliquido et al. 2023:**  
<https://arxiv.org/pdf/2303.15515.pdf>

# Pop. II BBHs: mass evolution



# detection rate

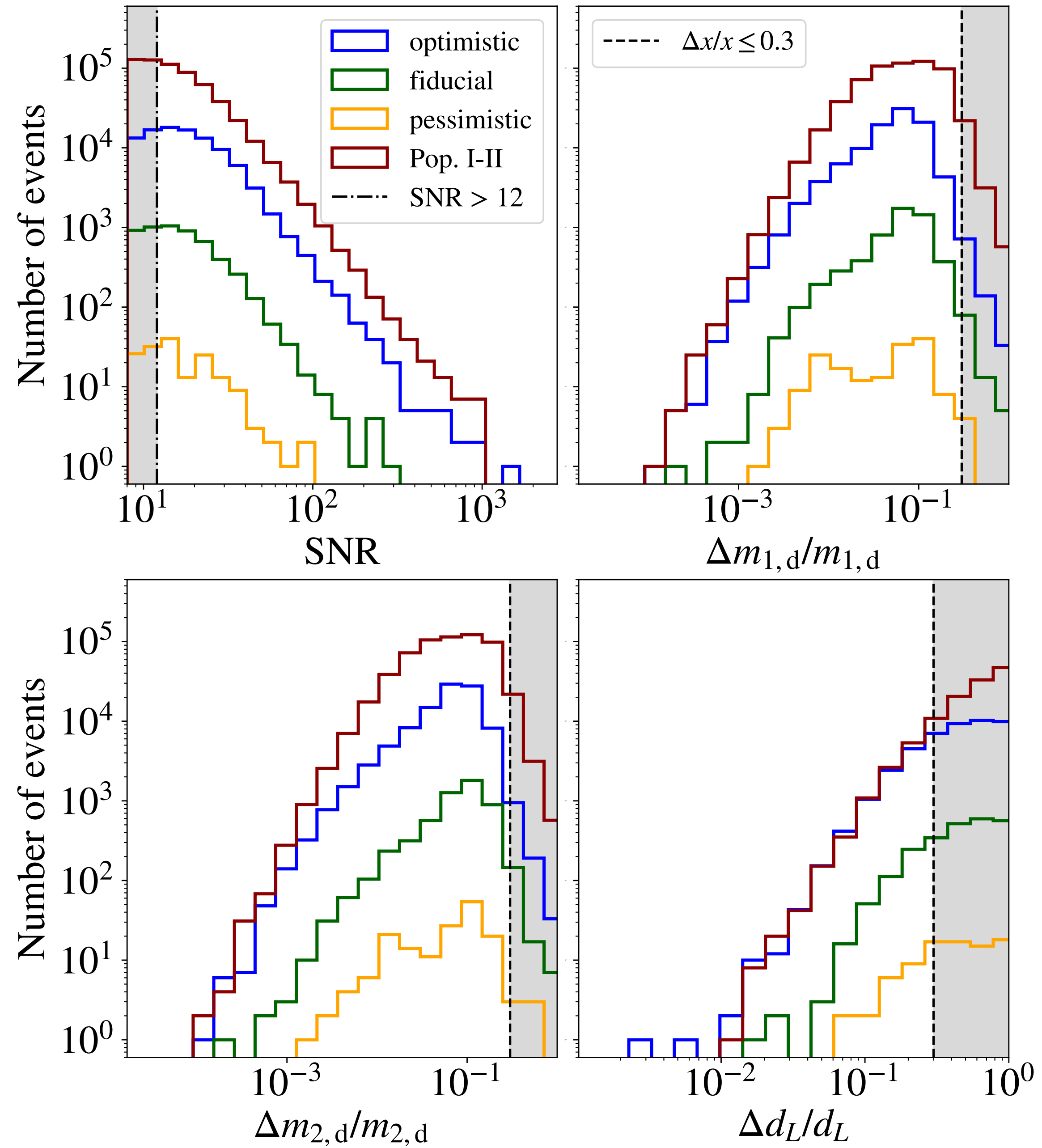
$$\mathcal{R}_{\text{det}} = \int \frac{d^2 \mathcal{R}(m_1, m_2, z)}{dm_1 dm_2} \frac{1}{(1+z)} \frac{dV_c}{dz} p_{\text{det}}(m_1, m_2, z) dm_1 dm_2 dz.$$

$$\frac{d^2 \mathcal{R}(m_1, m_2, z)}{dm_1 dm_2} = \mathcal{R}(z) p(m_1, m_2 | z).$$

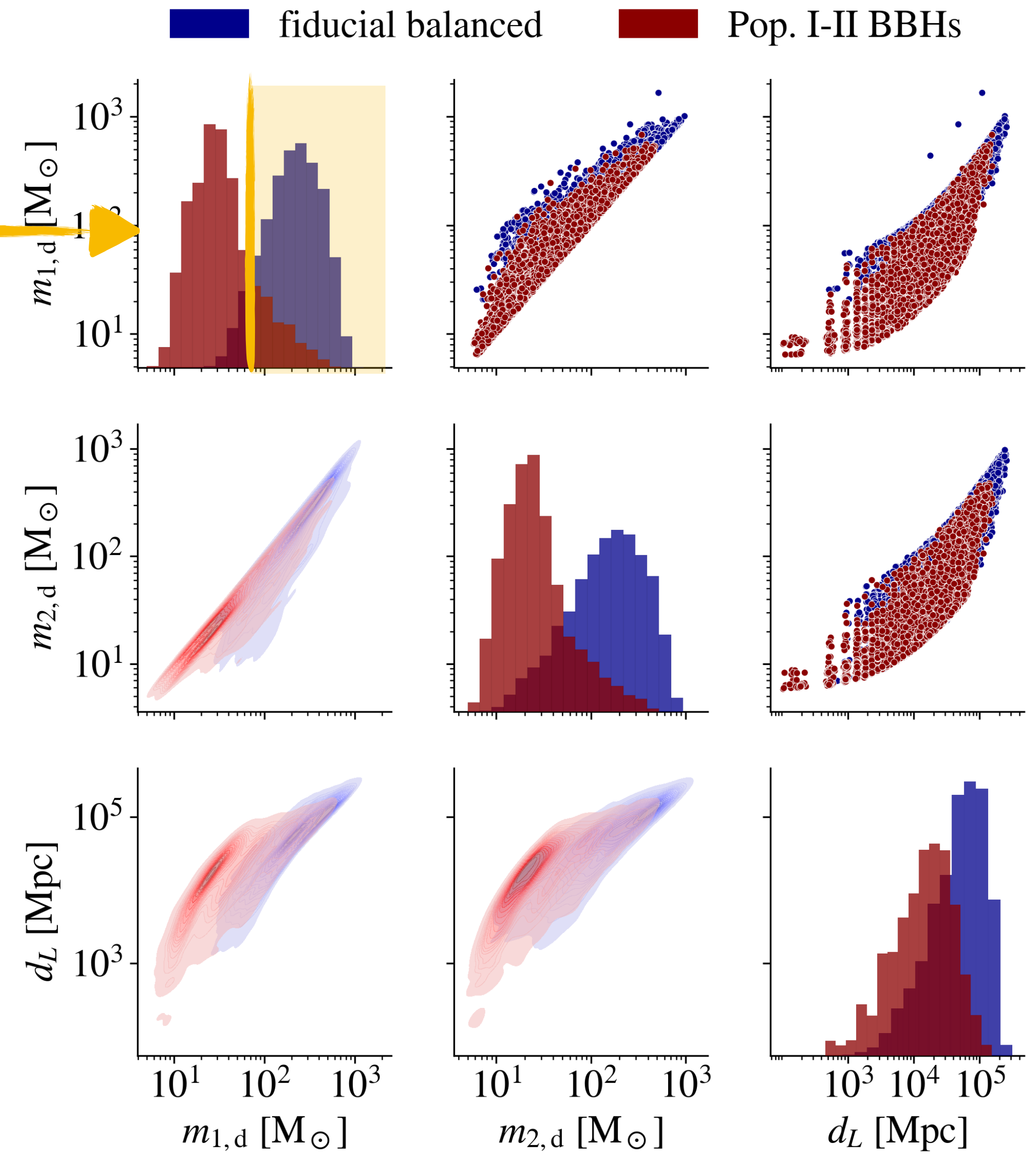
$$\rho = \rho_{\text{opt}} \sqrt{\omega_0^2 + \omega_1^2 + \omega_2^2}$$

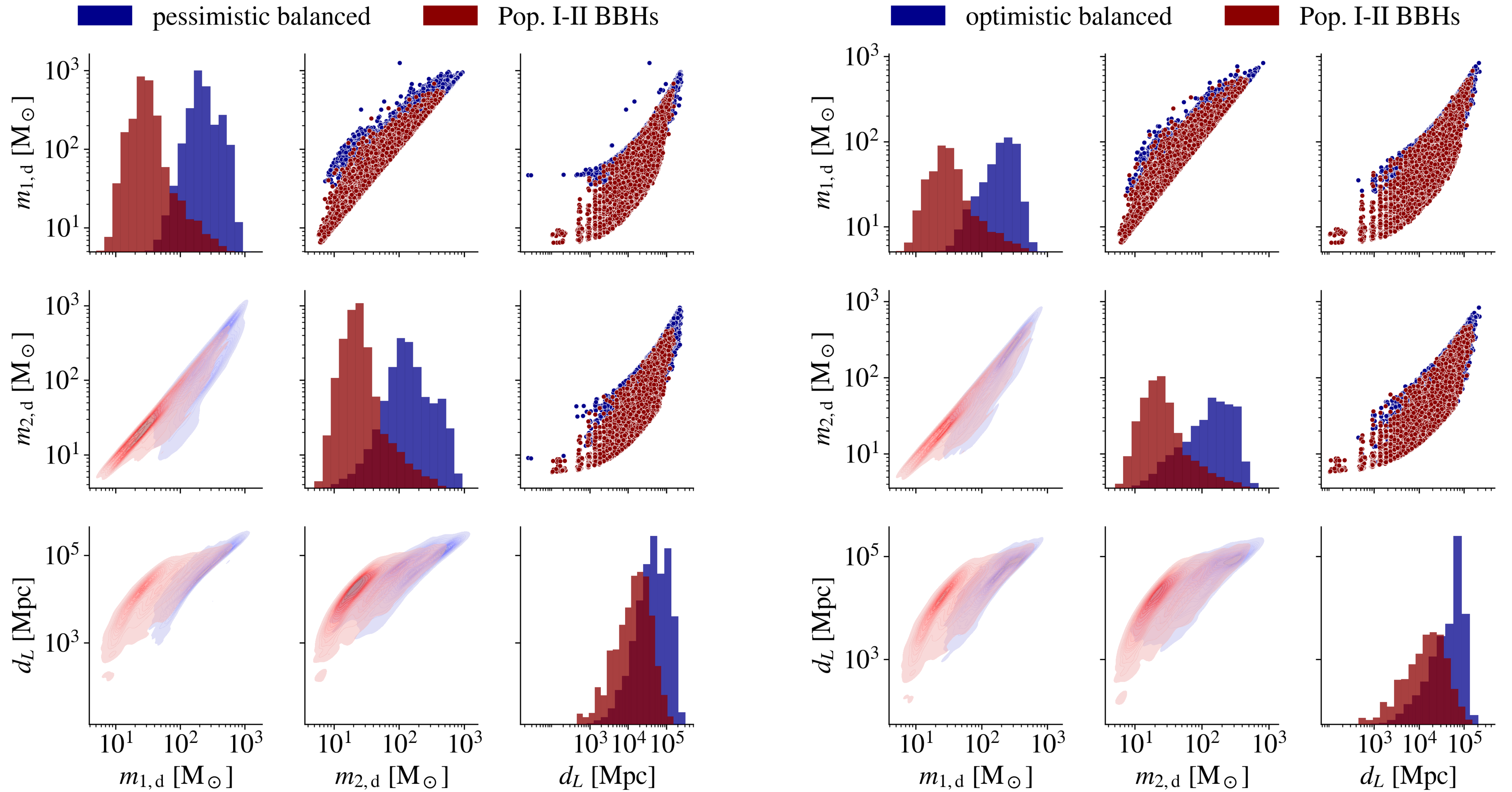
$$\rho_{\text{opt}}^2 = 4 \int_{f_{\text{low}}}^{f_{\text{high}}} df \frac{|\tilde{h}(f)|^2}{S_n(f)}$$

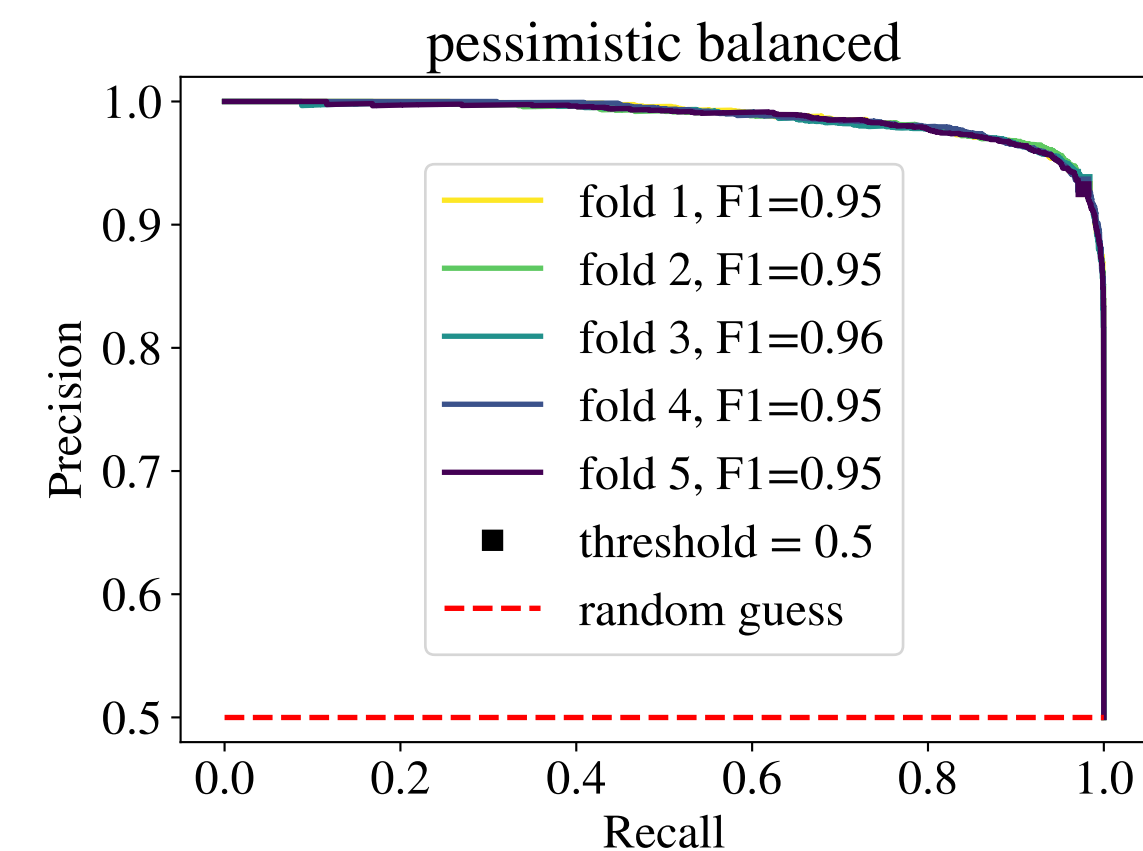
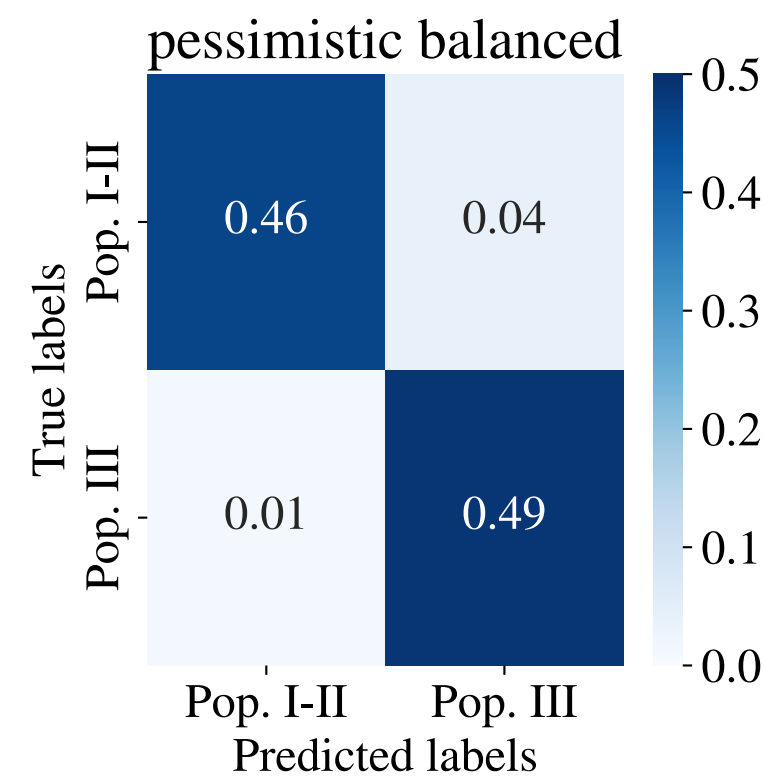
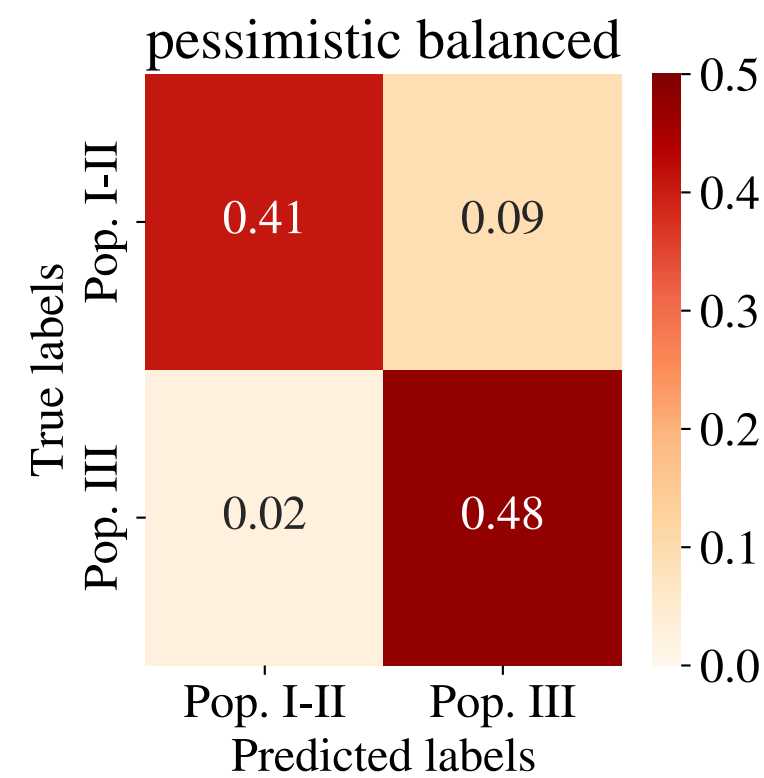
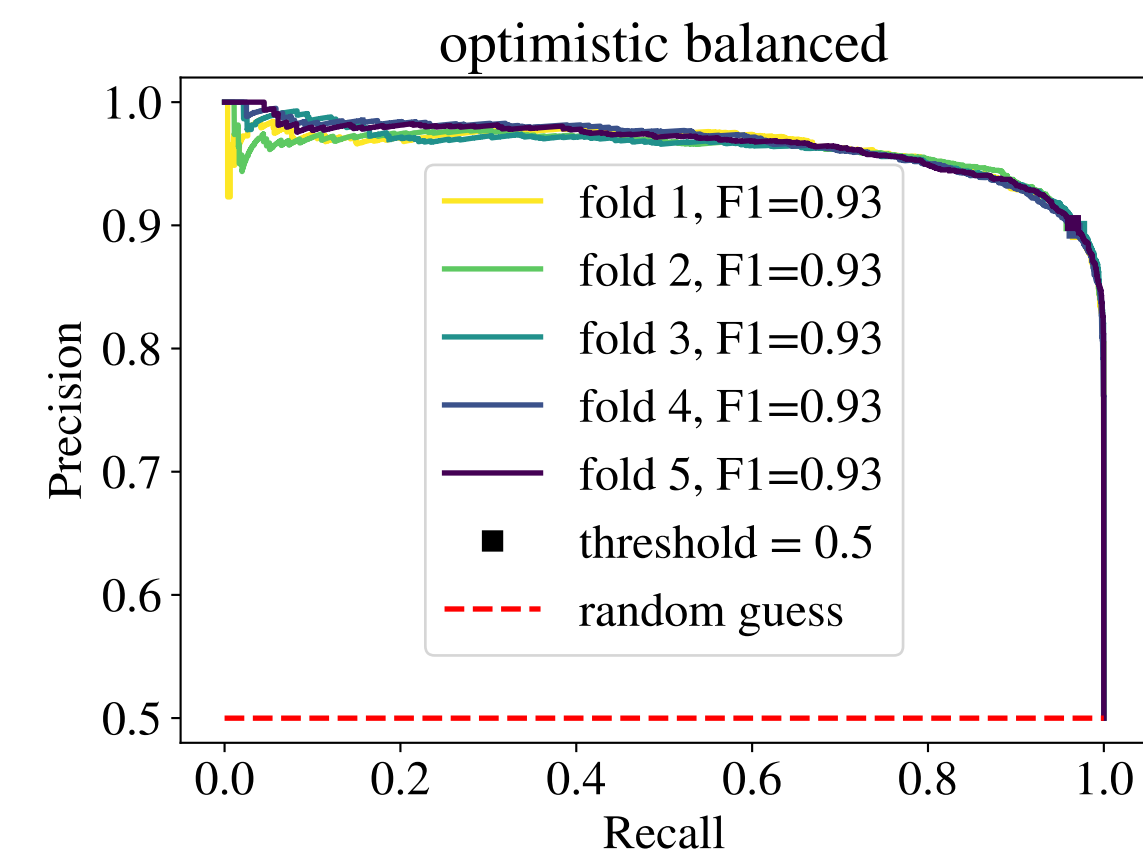
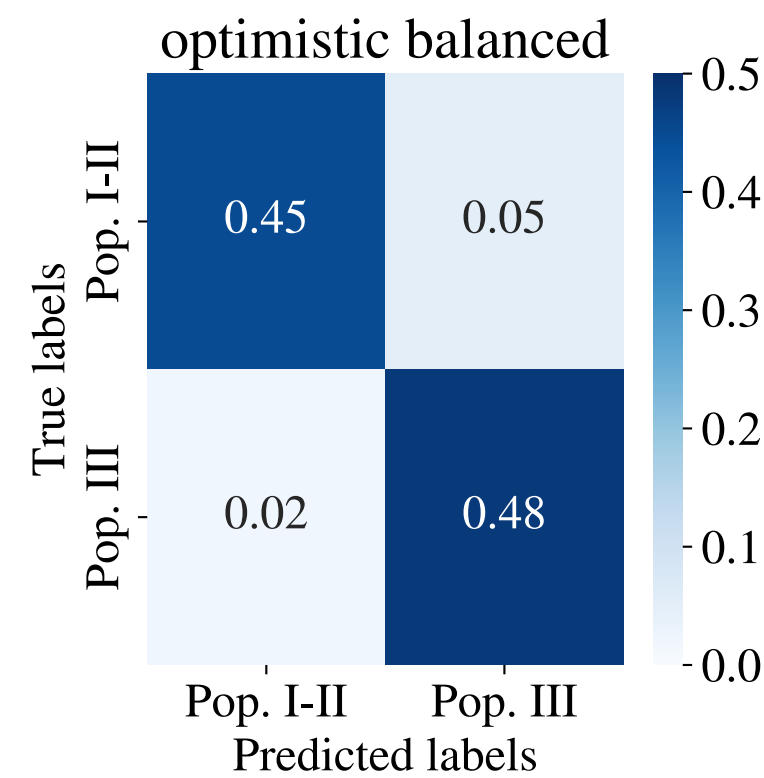
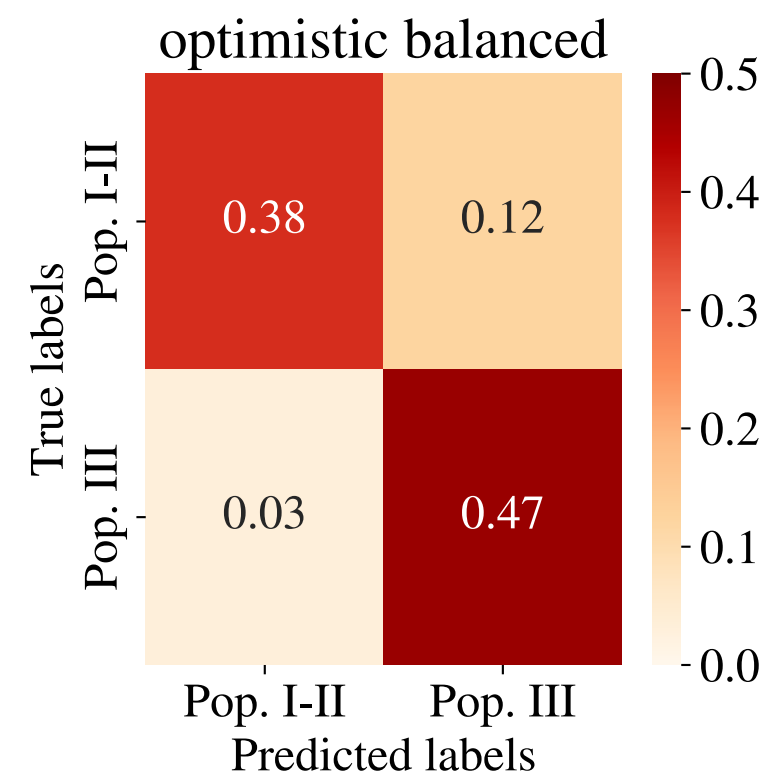
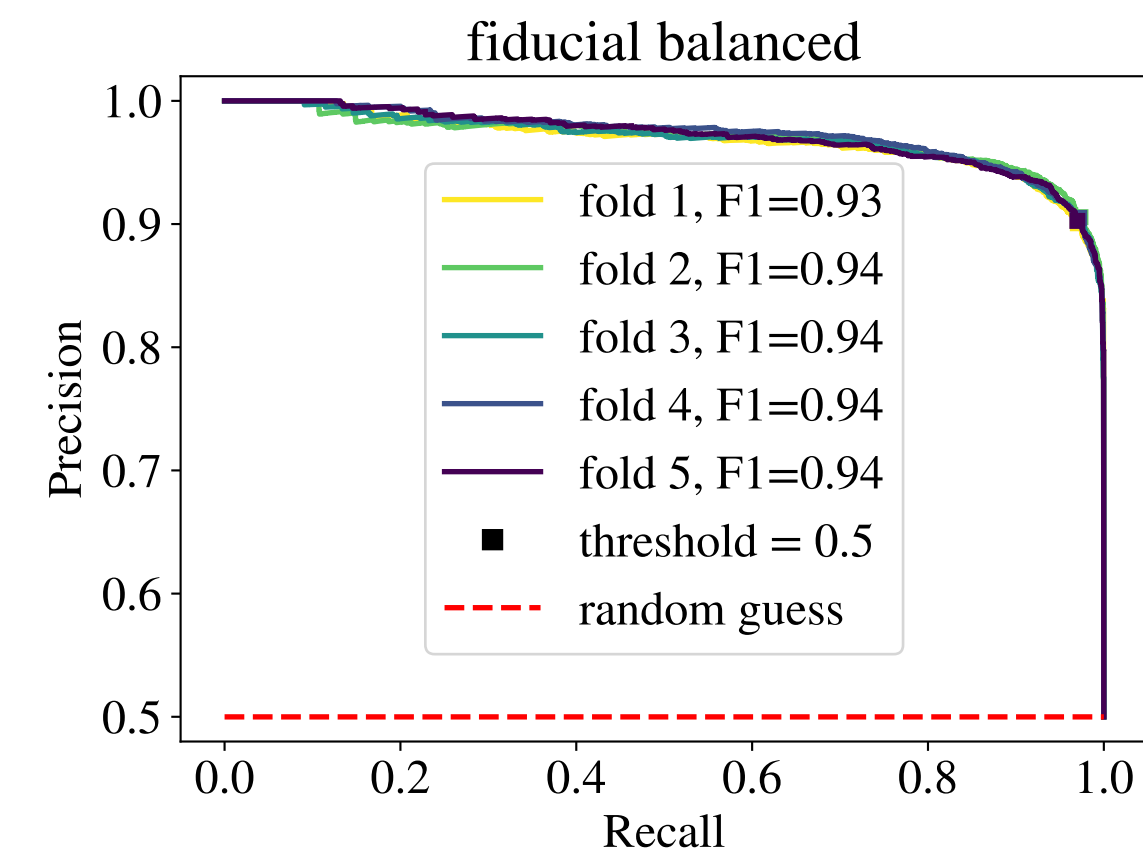
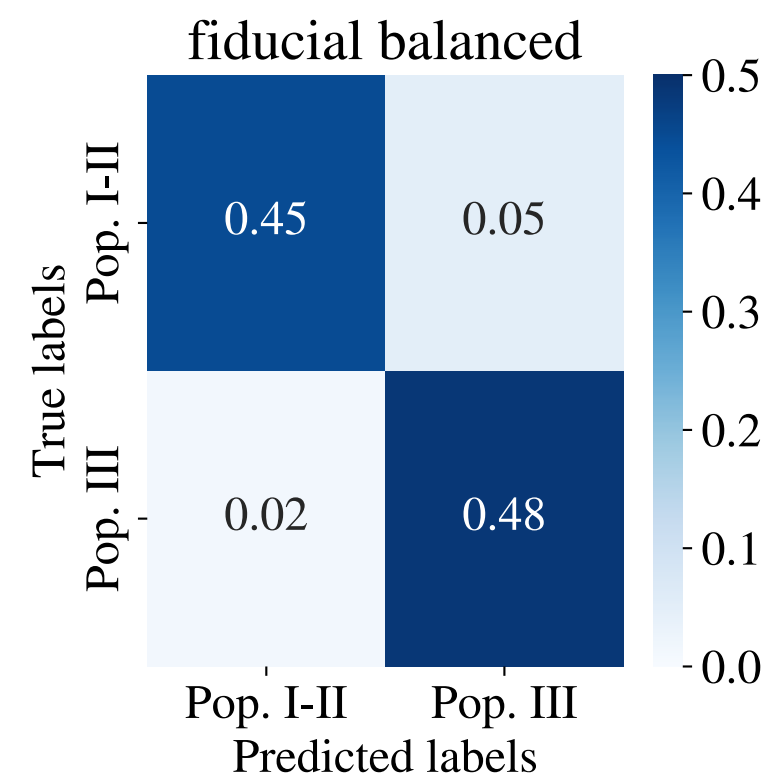
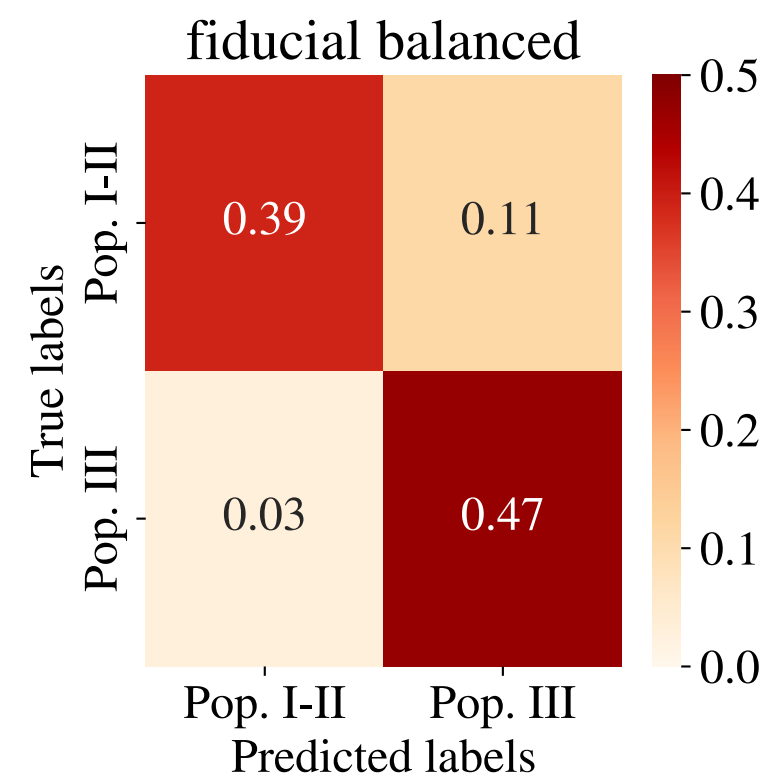




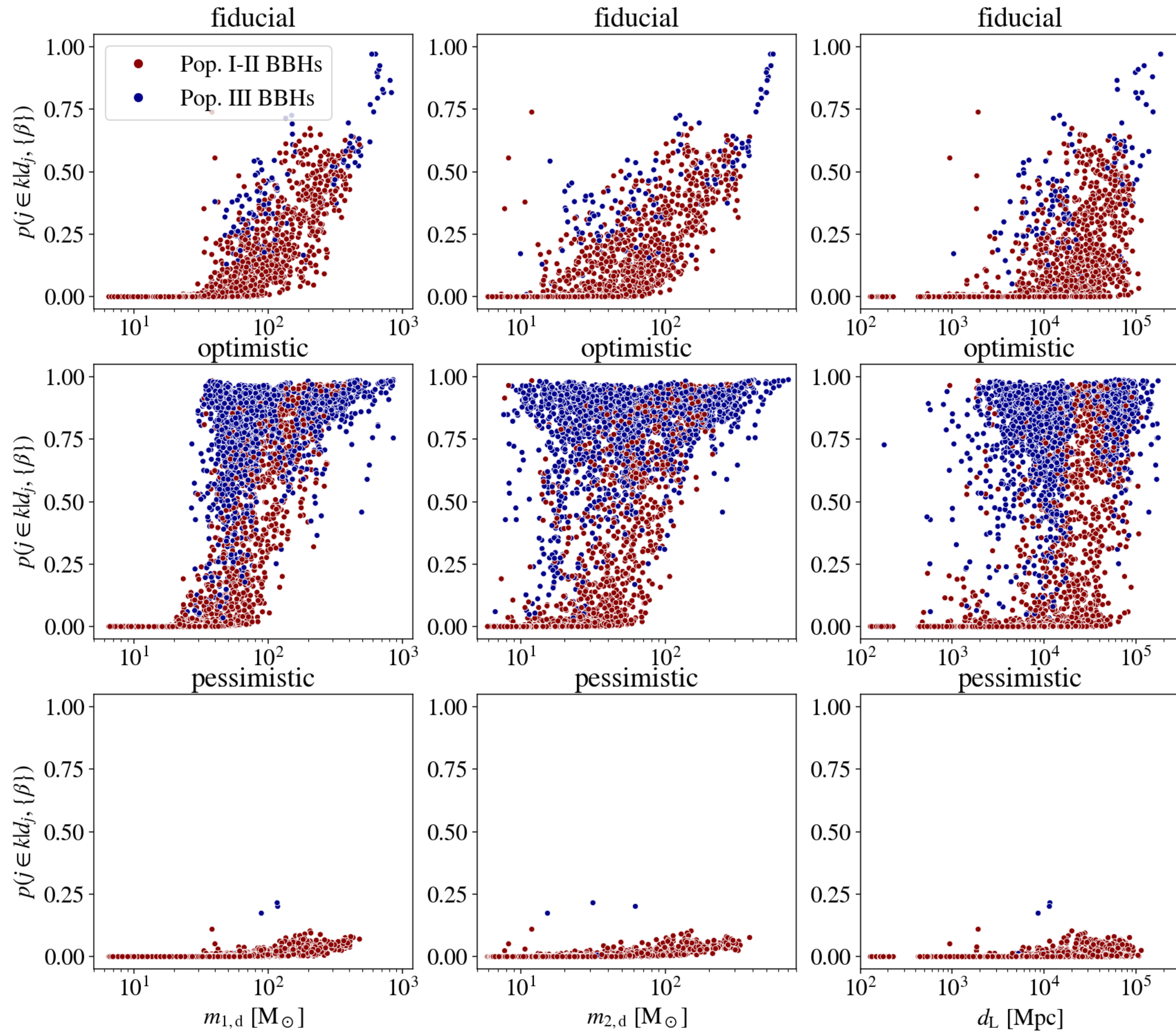
$$p(j \in k | x, d_j, \{\beta\}) = 1 \quad \text{if} \quad m_{1,d} \gtrsim 60 M_{\odot}$$











<b>Fiducial</b>						
Thr.	%TP	%TN	%FP	%FN	Precision	Recall
0.1	96	85	15	4	0.20	0.96
0.2	86	90	10	14	0.26	0.86
0.5	33	98	2	67	0.43	0.33
0.7	11	100	0	89	0.94	0.11
0.9	3	100	0	97	1.00	0.03

<b>Optimistic</b>						
Thr.	%TP	%TN	%FP	%FN	Precision	Recall
0.1	100	77	23	0	0.80	1.00
0.2	99	80	20	1	0.81	0.99
0.5	95	85	15	5	0.85	0.95
0.7	87	89	11	13	0.88	0.87
0.9	46	96	4	54	0.91	0.46

<b>Pessimistic</b>						
Thr.	%TP	%TN	%FP	%FN	Precision	Recall
0.1	50	100	0	50	0.60	0.50
0.2	33	100	0	67	1.00	0.33
0.5	0	100	0	100	0	0
0.7	0	100	0	100	0	0
0.9	0	100	0	100	0	0